

# Molecular Magnetism

## Theory & Experiment

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DIPC Community Seminar – March 16<sup>th</sup> 23



# Outline

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- Motivation
- Molecular Magnetism
  - Experimental characterisation
  - Theory and computation
- Single Molecule Magnets
  - What makes an SMMs and why do we focus on Dy(III)?
  - Molecular design: from static to dynamic properties
- Organic Radicals
  - How to trick metal-free molecules into not forming bonds: topology
  - Impact of structural flexibility
- What's next?

# Outline

- Motivation



- Molecular Magnetism

- Experimental characterisation
- Theory and computation

~ 20 minutes

- Single Molecule Magnets

- What makes an SMMs and why do we focus on Dy(III)?
- Molecular design: from static to dynamic properties

~ 10 minutes

- Organic Radicals

- How to trick metal-free molecules into not forming bonds: topology
- Impact of structural flexibility

~ 10 minutes

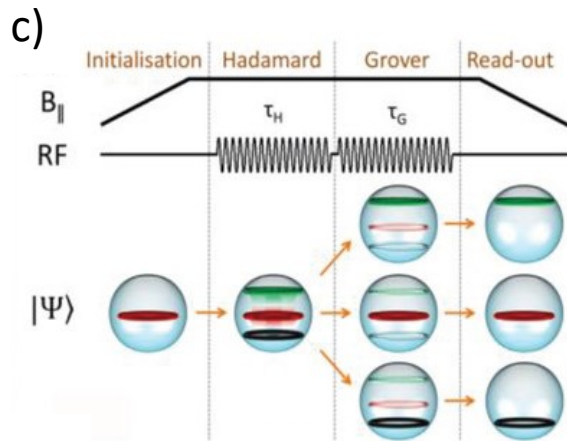
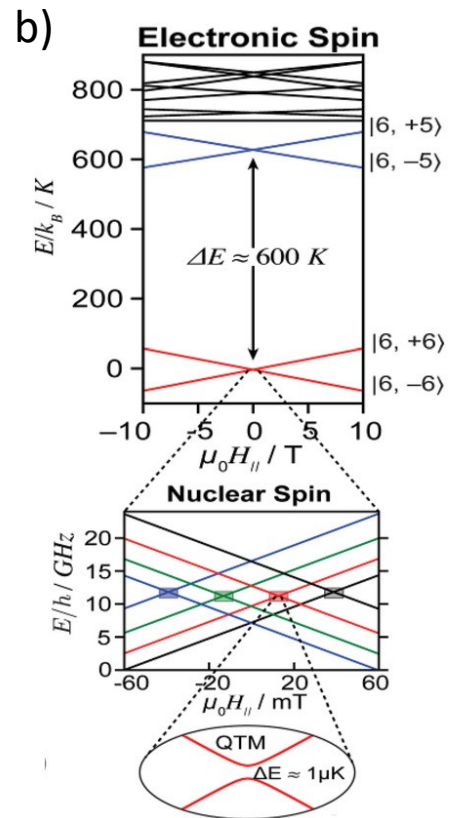
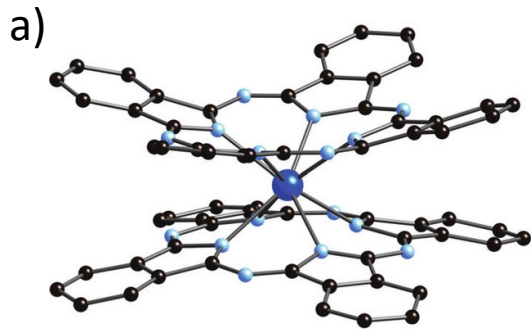
- What's next?

~ 5 minutes

# Motivation

- *Quantum information*

Grover algorithm in a  $^{159}\text{Tb(III)Pc}_2$   
**qudit** single molecule transistor



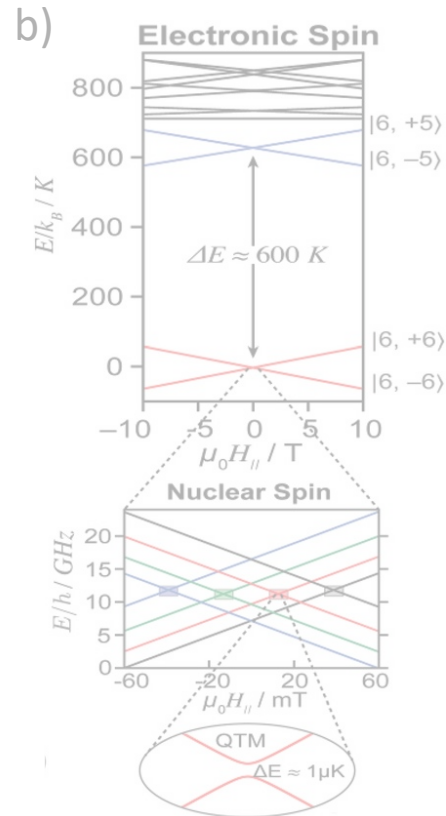
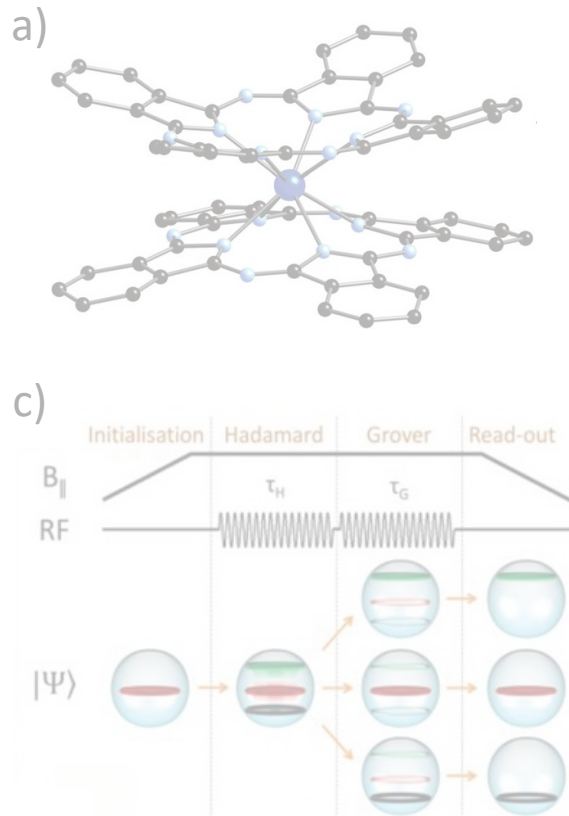
*Chem. Soc. Rev.*, **2012**, *41*, 7464-7478

*Phys. Rev. Lett.* **2017**, *119*, 187702 *Adv. Mater.* **2019**, *31*, 1806687

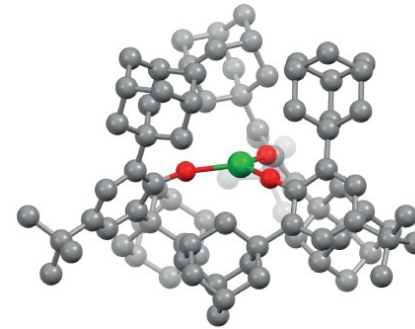
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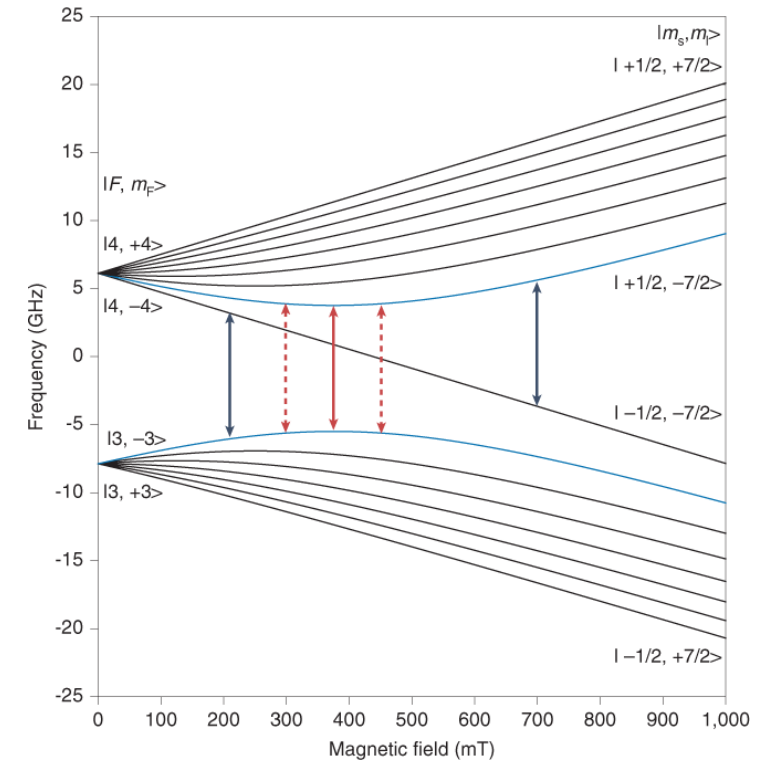


Improving **quantum coherences** by operating at clock transitions (CT) in a  $\text{Lu(II) } 5d^1$



$$dE/dB \propto M$$

At CT, system is protected from magnetic noise (dipolar fields)



$$F = I + S$$

$$E \propto B_0 m_F$$

$$\mu_B \vec{B} \vec{g} \vec{S}$$

$$E \propto B_0 m_S$$

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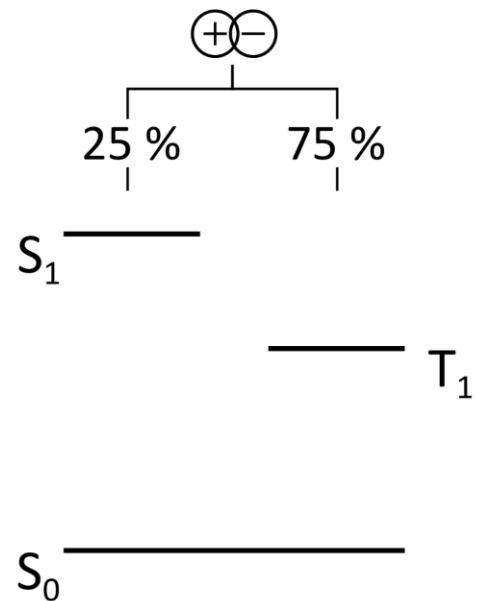
Nat. Chem. 2022, 14, 361-362 Nat. Chem. 2022, 14, 392-397

# Motivation

- *OLEDs*

**Spin-statistics** limit Quantum Efficiency

Conventional fluorescence



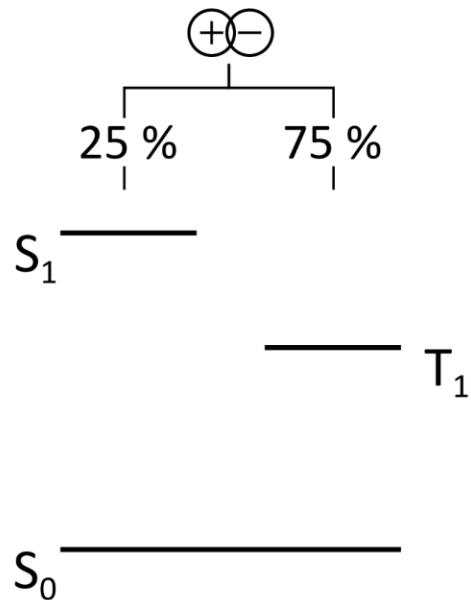
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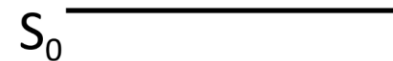
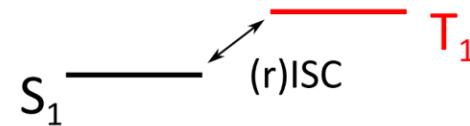
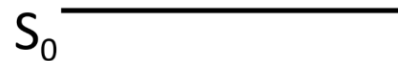
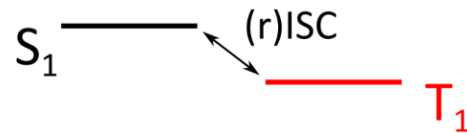
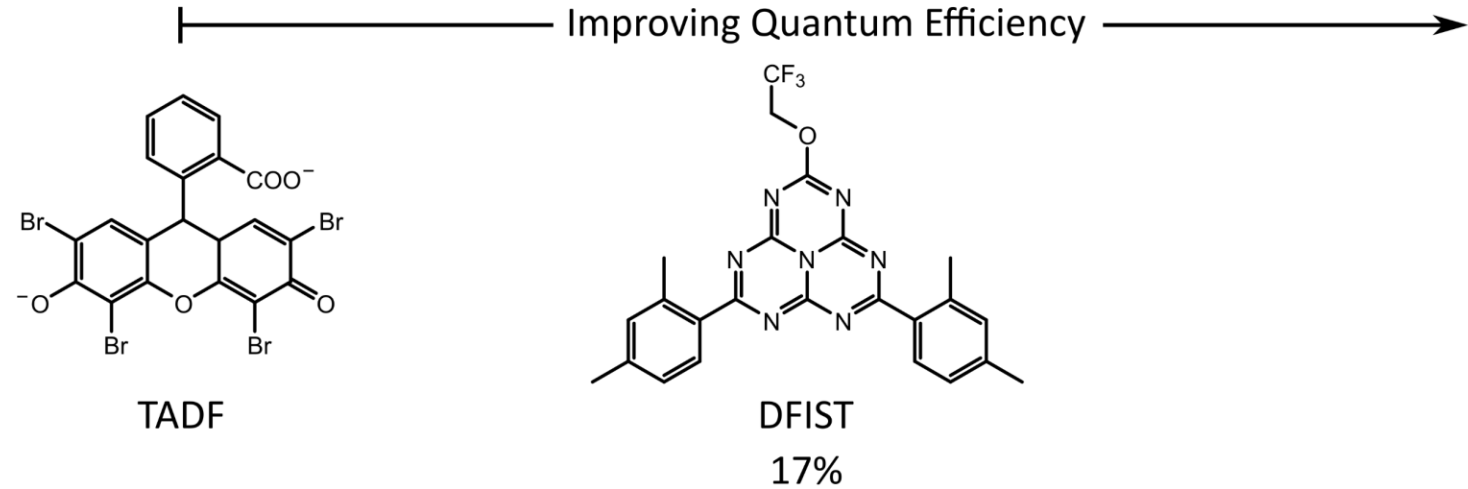
Spin-statistics limit Quantum Efficiency

Involve  $T_1$  ( $\Delta E_{ST}$  & SOC) to improve QE

Conventional fluorescence



*Chem. Soc. Rev.*, **2017**, 46, 915-1016



*Nature*, **2022**, 609, 502-506

# Motivation

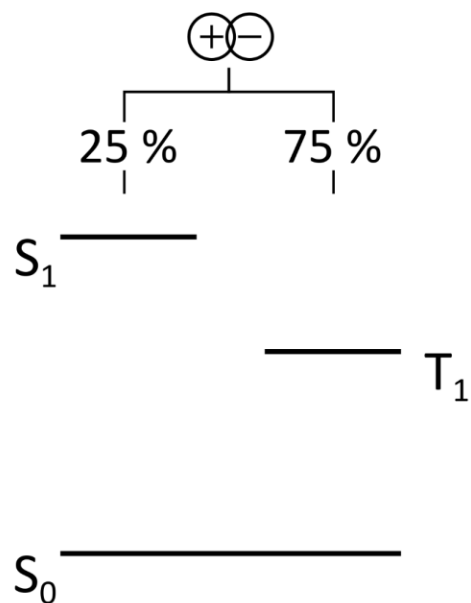
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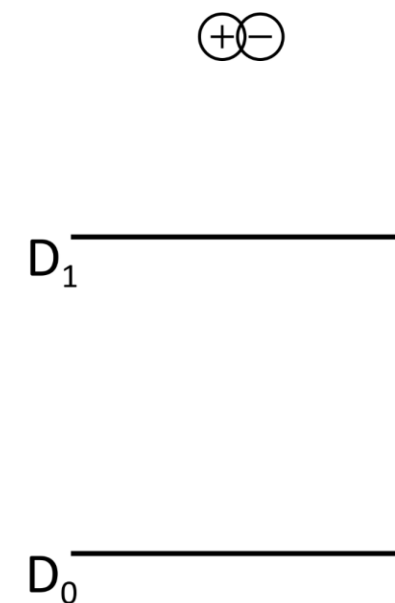
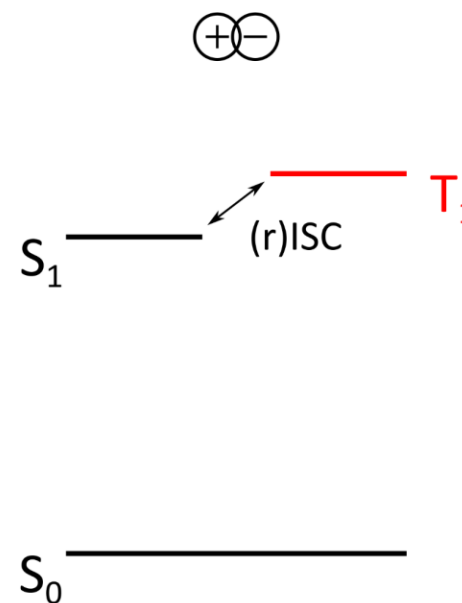
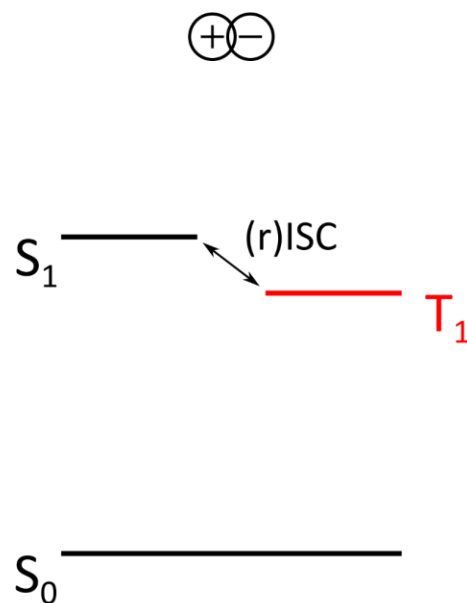
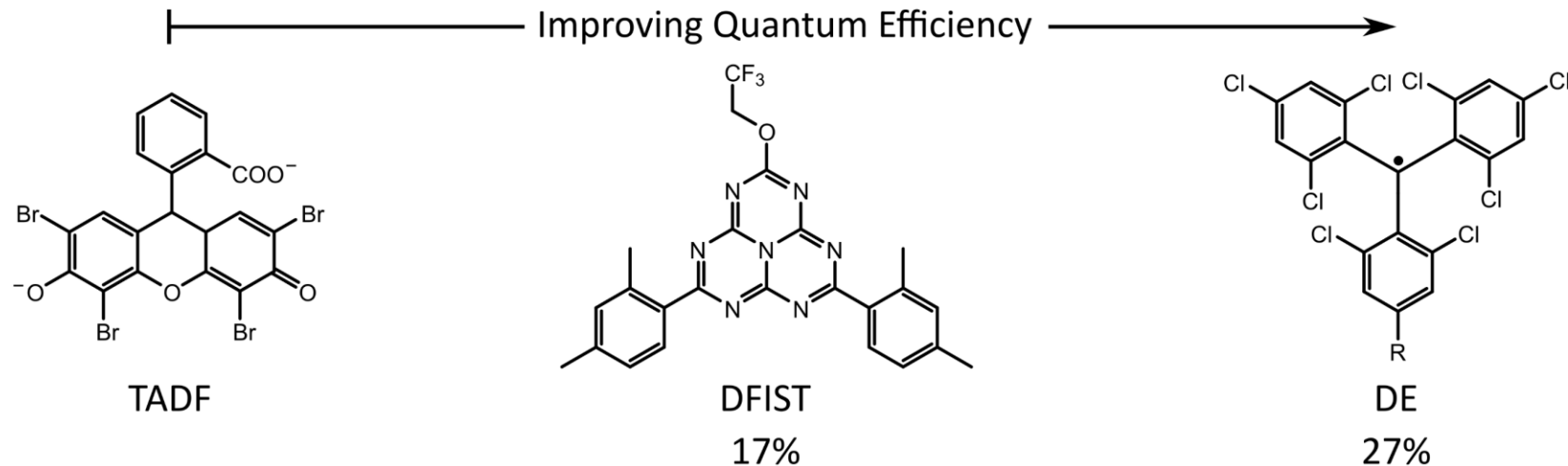
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Spin doublet avoids them altogether

Conventional fluorescence



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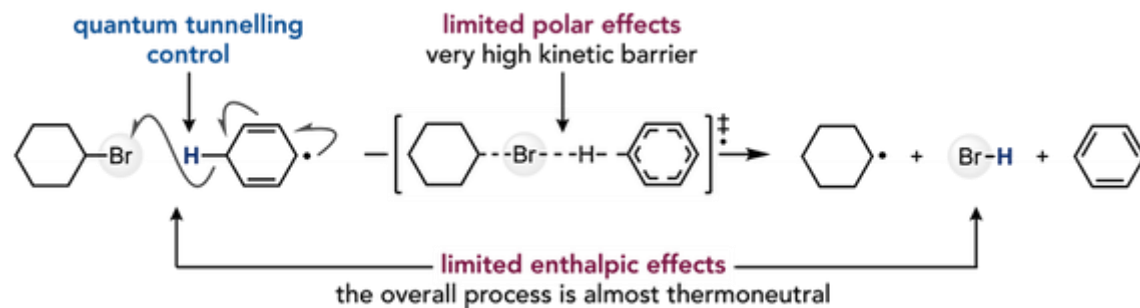
# Motivation

- Synthetic chemistry

Radical-mediated routes to C-H bond formation

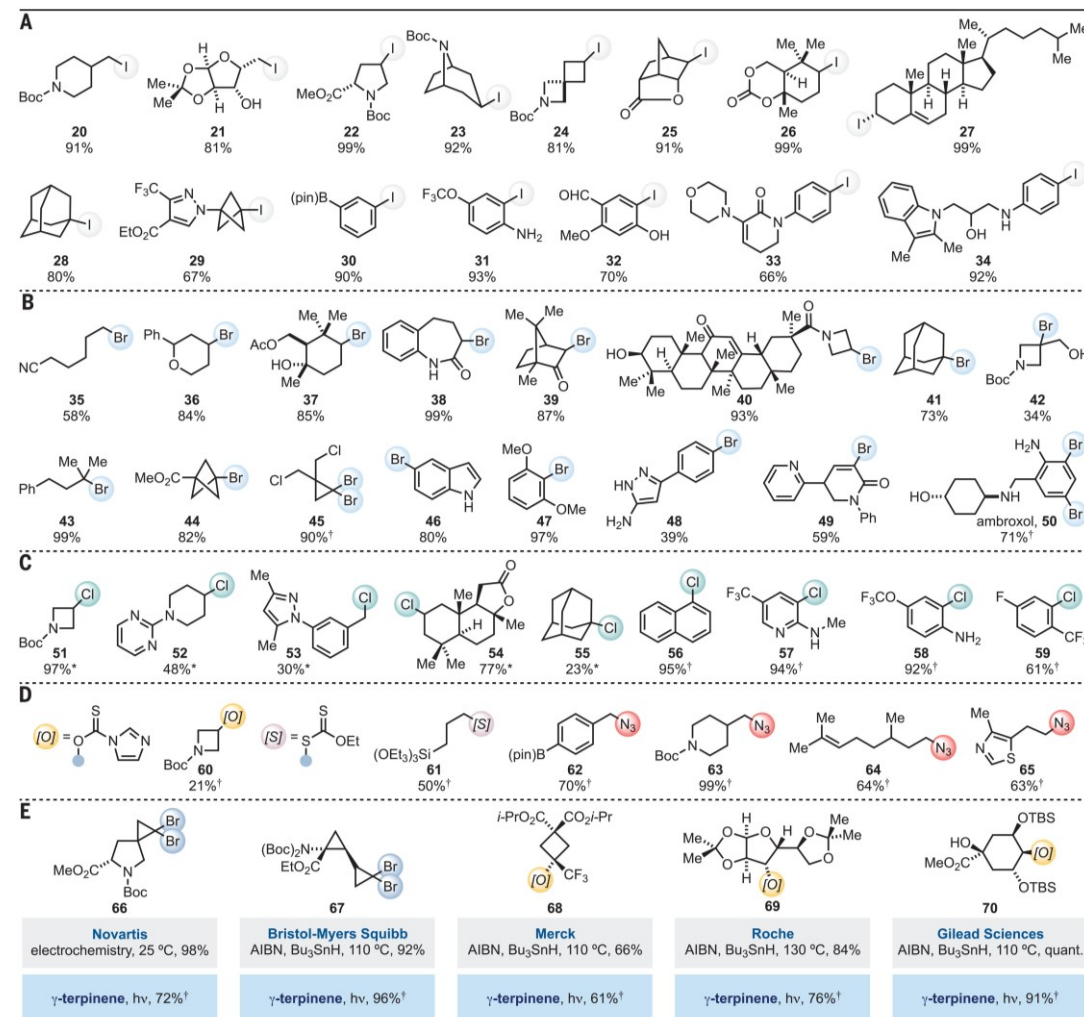
High yields and “mild” conditions

By-product: volatile aromatic in place of tin or silicon waste



Science, 2022, 377, 1323–1328

Massive substrate scope



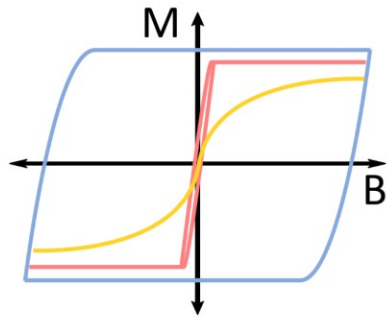
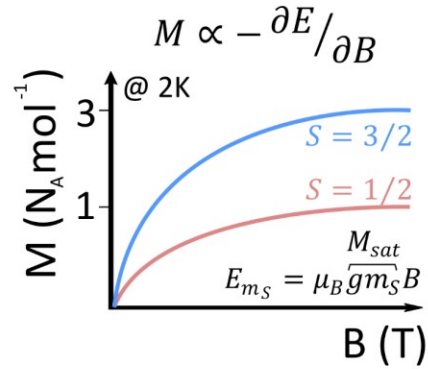
# Molecular Magnetism

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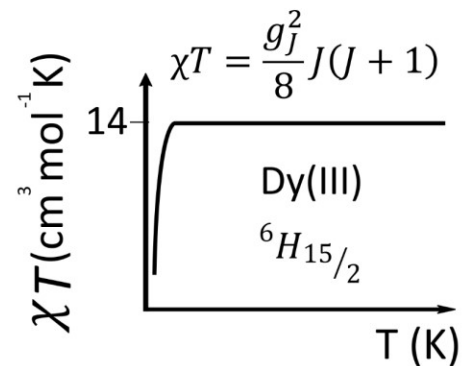
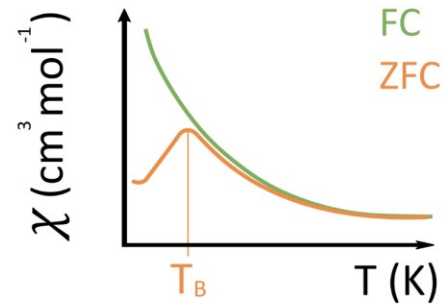
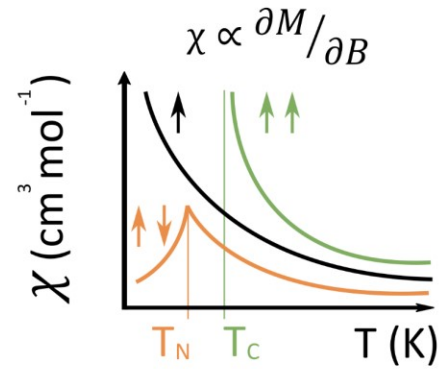
- *Experimental characterisation –*

# Molecular Magnetism

- Experimental characterisation – DC magnetometry measures moment varying field and temperature

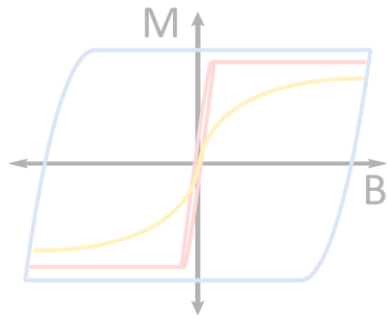
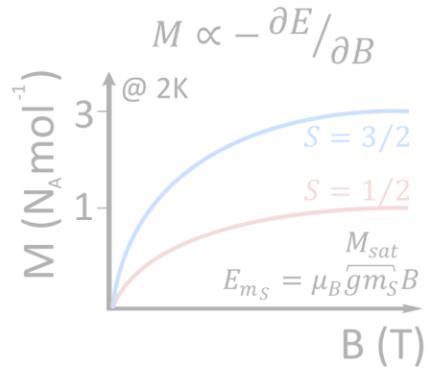


Access to **static & dynamic** properties of **spin-only** and **spin-orbit** coupled systems

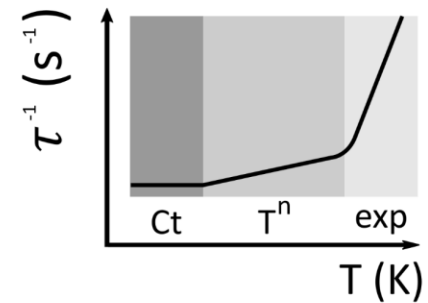
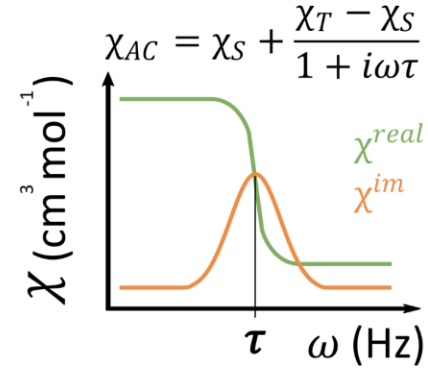
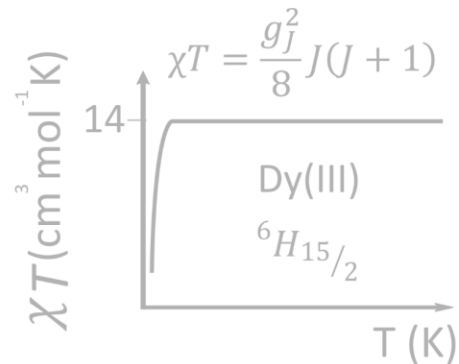
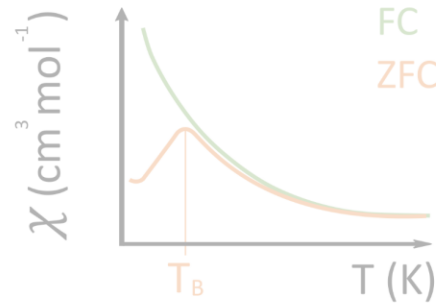
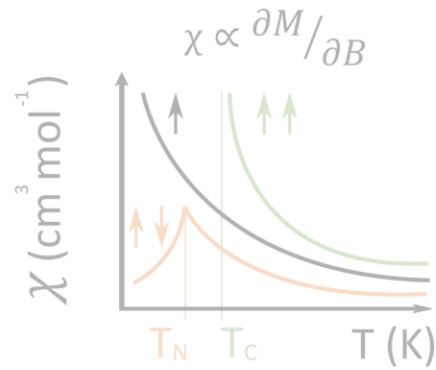


# Molecular Magnetism

- *Experimental characterisation* – **AC** magnetometry measures moment varying field, temperature and AC frequency



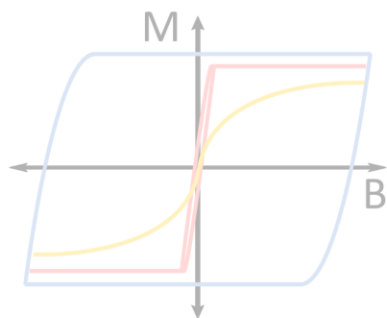
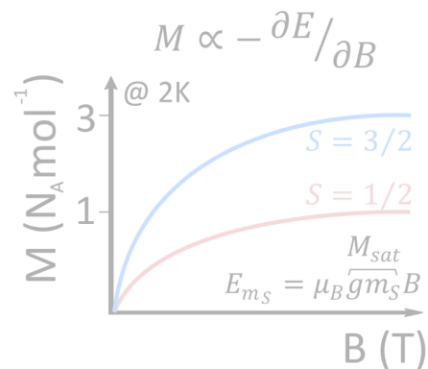
Access to **static** & **dynamic** properties of **spin-only** and **spin-orbit** coupled systems



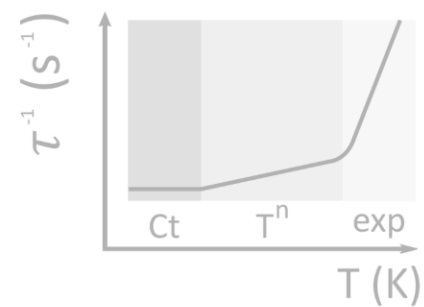
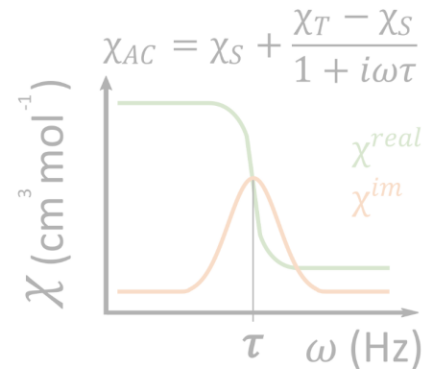
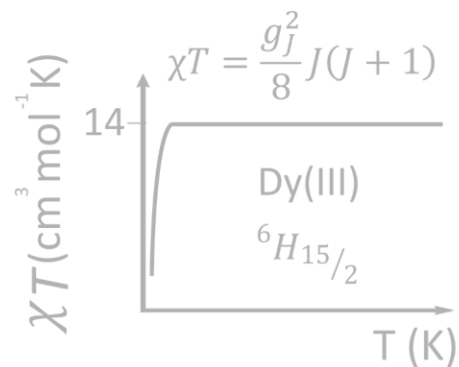
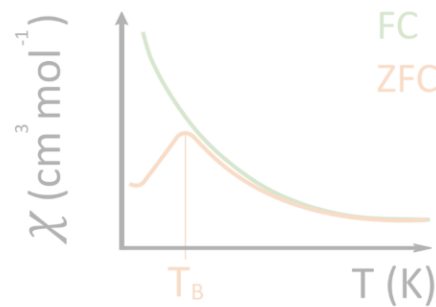
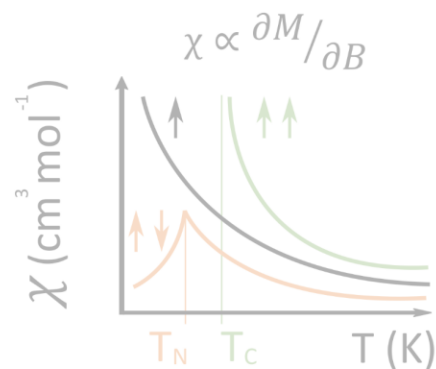
Access to relaxation of magnetisation  
**characteristic time ( $\tau$ )**

# Molecular Magnetism

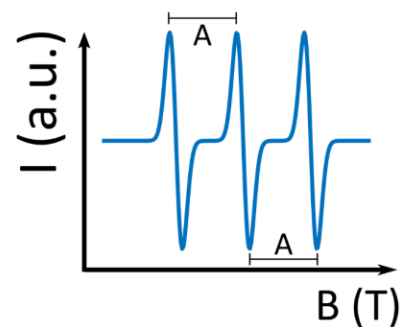
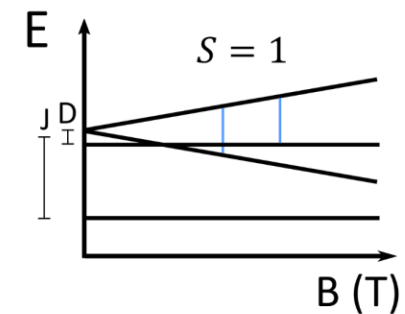
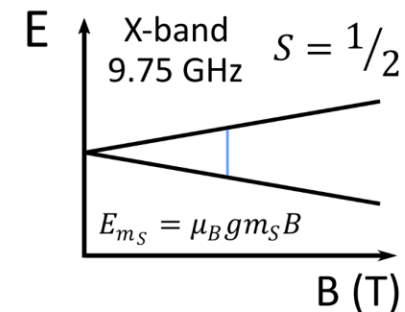
- Experimental characterisation – EPR measures spin transitions varying field and temperature



Access to **static & dynamic** properties of **spin-only** and **spin-orbit** coupled systems

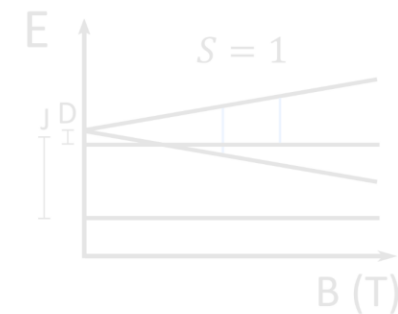
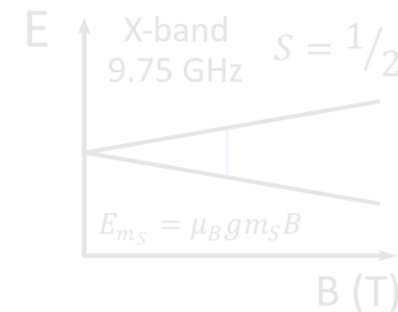
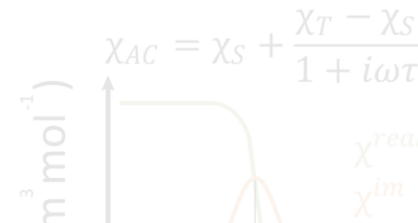
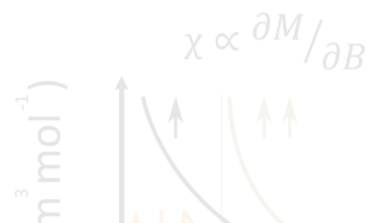
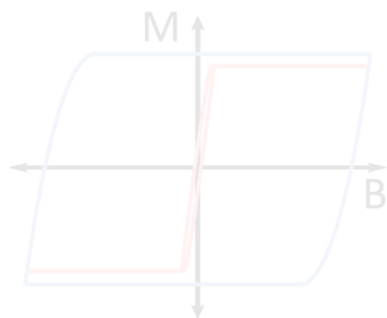
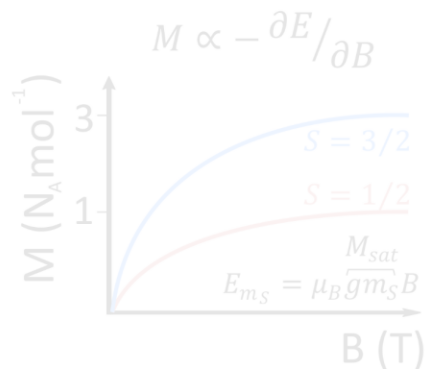


Access to relaxation of magnetisation **characteristic time ( $\tau$ )**



# Molecular Magnetism

- *Experimental characterisation* – electronic structure interrogated with complementary methods



- A physically meaningful model spin Hamiltonian is formulated for the system
- Data is fitted (global) to validate, refine or discard the model
- Initial guess values are obtained from ab initio methods
- PHI<sup>1</sup> or EasySpin<sup>2</sup> are customarily used

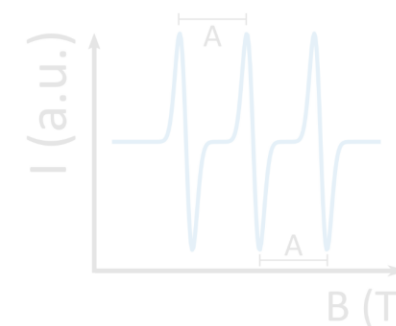
1 *J. Comp. Chem.*, **2013**, 34, 1164–1175

2 *J. Magn. Reson.* **2006**, 178(1), 42-55

Access to static & dynamic properties of spin-only and spin-orbit coupled systems



magnetisation characteristic time ( $\tau$ )



# Molecular Magnetism

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- *Theory and computation* –

# Molecular Magnetism

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- *Theory and computation* – **Geometry optimisation** with Density Functional Theory (DFT)
  - Both in gas-phase (*Gaussian*) and periodic conditions (*Crystal, VASP, Phonopy*)
  - For spin-only systems:
    - Hybrid, GGA, meta-GGA, range-corrected functionals & standard basis sets
  - For spin-orbit coupled systems:
    - Exchange-correlation functionals & ECPs – diamagnetic analogue with actual isotopic mass for normal modes



# Molecular Magnetism

- *Theory and computation* – **Electronic structure** with DFT & multiconfigurational methods (CASSCF, PT2, MRCI)

➤ For spin-only systems:

- $\Delta E$  between spin-adapted (*OpenMolcas*) states describe model spin Hamiltonian parameters ( $J$ , ZFS,  $t/U$ , ...)
- Mapping approaches to employ broken symmetry solutions (*Gaussian*)
- Standard basis sets

Exact

$$\hat{H}_{elec} = - \sum_{i=1}^N \frac{1}{2} \nabla_i^2 - \sum_{i=1}^N \sum_{A=1}^M \frac{Z_A}{r_{iA}} + \sum_{i=1}^N \sum_{j>1}^N \frac{1}{r_{ij}}$$

Non-relativistic & time-independent

Spin-adapted

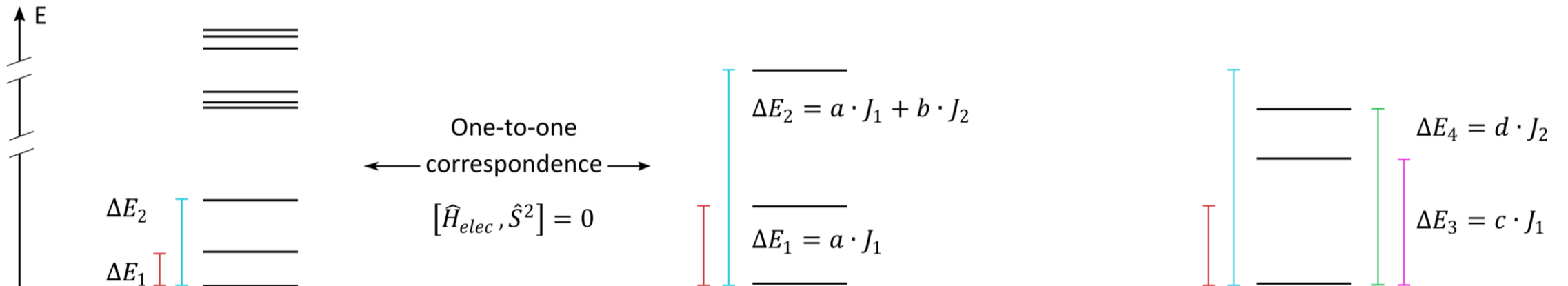
$$\hat{H}^{HDVV} = - \sum_{\langle i,j \rangle} J_{ij} \hat{S}_i \cdot \hat{S}_j$$

$$|\Psi\rangle_{MC} = \sum_I c_I |\phi_I\rangle$$

broken-symmetry

$$\hat{H}^{Ising} = - \sum_{\langle i,j \rangle} J_{ij} \hat{S}_i^z \cdot \hat{S}_j^z$$

$$|\Psi\rangle_{BS} = |\phi\rangle$$



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➤ For spin-orbit coupled systems (*OpenMolcas*) :

- Relativistic corrections: scalar 2<sup>nd</sup>-order Douglas-Kroll-Hess and ANO-RCC basis set
- SOC: AMFI and State-Interaction between appropriate number of states for each spin-multiplicity
- Parametrisation of ab initio results by means of Crystal Field Theory

$$\hat{H}_{CF} = \sum_{k=2,4,6} \sum_{q=-k}^k B_k^q \theta_k \hat{O}_k^q$$

$B_k^q$ : Crystal Field Parameters (CFPs)

$\theta_k$ : Operator equivalent factors

$\hat{O}_k^q$ : Stevens operators (LC of  $\hat{L}_Z, \hat{L}_+, \hat{L}_-$ )

# Molecular Magnetism

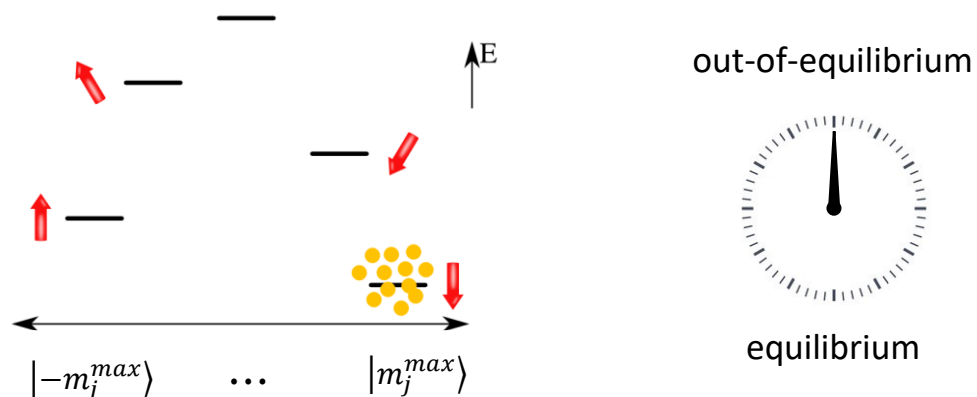
- *Theory and computation* – **Spin dynamics** with master matrix & 1<sup>st</sup> order TDPT

- Description of spin-phonon coupling:

- Gas-phase molecular normal modes
- Transition rates are calculated with Fermi's golden rule
- Orbach process

- Dynamics:

- Master matrix formalism
- Solution yields how long it takes to reach equilibrium  $\tau$
- Solve at different temperatures for direct comparison to experimental relaxation profiles
- Identify main deactivation pathways and *hope* to affect them with molecular design



# Single Molecule Magnets

---

- *What makes an SMM and why Dy(III)?*

**Slow relaxation of magnetisation** → roughly  $10^5 < \tau(s) < 10^2$  at 2 K,  $\tau(s) < 10^{-4}$  as hot as possible

- Key ingredient is magnetic anisotropy, achieved with large spin-orbit coupling (SOC)

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*Example:* Dy(III)  $4f^9$  electronic configuration.  $4f$  orbitals are shielded by filled atomic shells → unquenched  $L$

$$\begin{array}{ccccccc} \uparrow\downarrow & \uparrow\downarrow & \uparrow & \uparrow & \uparrow & \uparrow & \uparrow \\ m_L: & -3 & -2 & -1 & 0 & 1 & 2 & 3 & L = \sum m_L = 5 \\ m_S: & 0 & 0 & 1/2 & 1/2 & 1/2 & 1/2 & 1/2 & S = \sum m_S = 5/2 \end{array}$$

Russell-Saunders coupling scheme:  $|L - S| \leq J \leq |L + S|$

$$5/2 \leq J \leq 15/2$$

How are they distributed?

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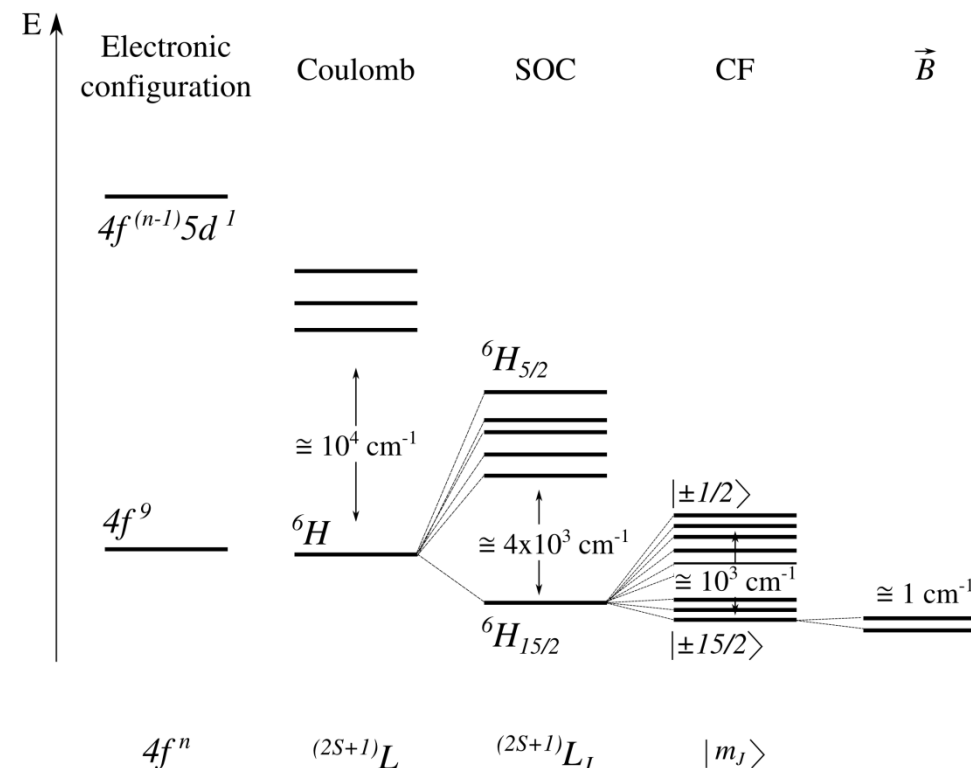
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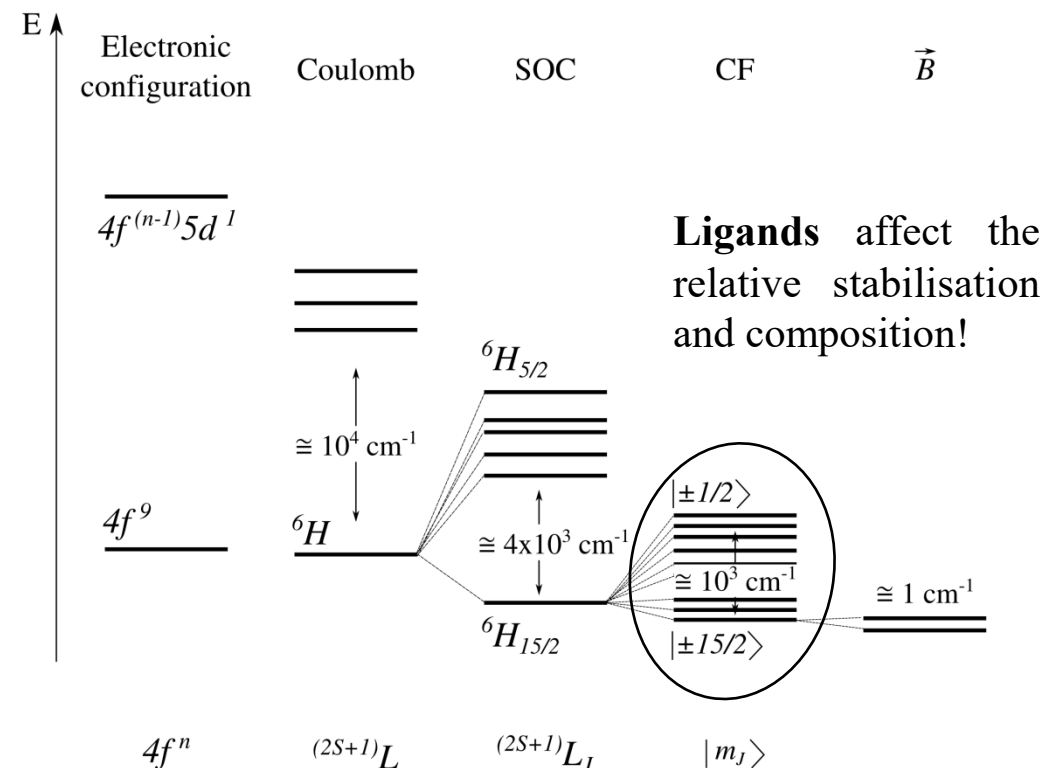
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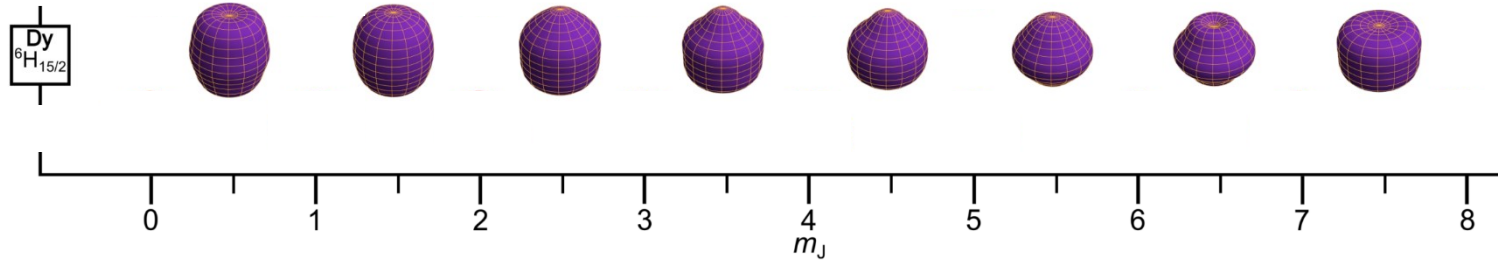


# Single Molecule Magnets

- *Molecular design: from static to dynamic properties*

**Axial** ligand fields are key

- States with largest  $m_J$  are stabilised; states with smallest  $m_J$  are destabilised –  $\Delta E$  informs of  $U^{eff}$



**Fig.** Free-ion charge density plots for the  $m_J$  states of the ground Hund's rule term of  $Dy(III)$  as derived by Sievers.<sup>1</sup> Figure credit: J. G. C. Kragsskow.

<sup>1</sup> J. Sievers, Zeitschrift für Physik B, Condensed Matter **1982**, 45, 289-296.

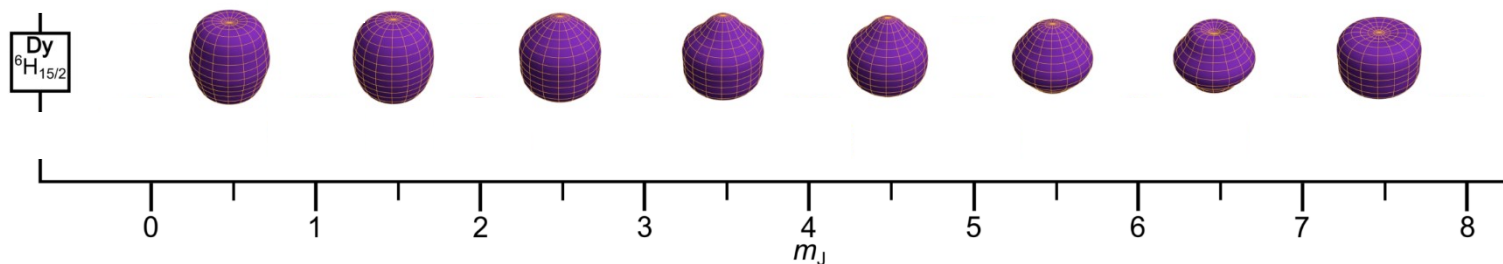


# Single Molecule Magnets

- *Molecular design: from static to dynamic properties*

**Axial** ligand fields are key

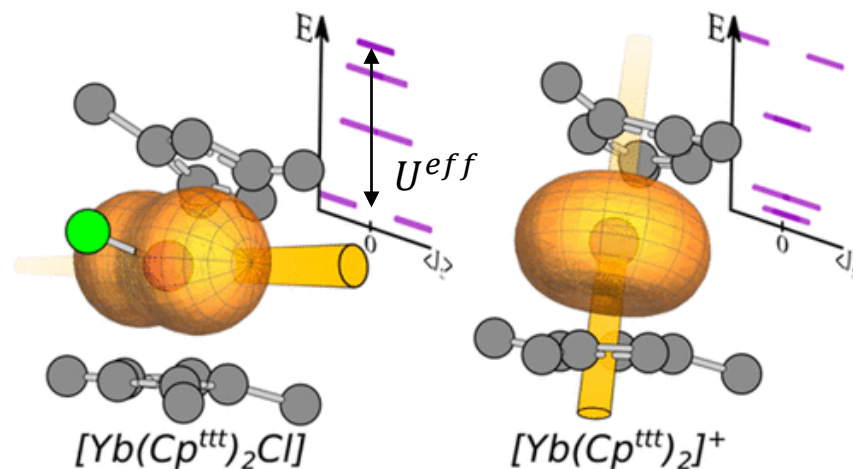
- States with largest  $m_J$  are stabilised; states with smallest  $m_J$  are destabilised –  $\Delta E$  informs of  $U^{eff}$



**Fig.** Free-ion charge density plots for the  $m_J$  states of the ground Hund's rule term of Dy(III) as derived by Sievers.<sup>1</sup> Figure credit: J. G. C. Kragoskow.

1 J. Sievers, Zeitschrift für Physik B, Condensed Matter **1982**, 45, 289-296.

- $U^{eff}$  is a key indicator of a SMM performance (static) – How easy it is to relax over the barrier



Note that for Yb(III), stabilisation is reversed wrt Dy(III)

C. A. P. Goodwin, D. Reta, F. Ortu, N. F. Chilton, D. P. Mills, *JACS* **2017**, 139, 51, 18714–18724

# Single Molecule Magnets

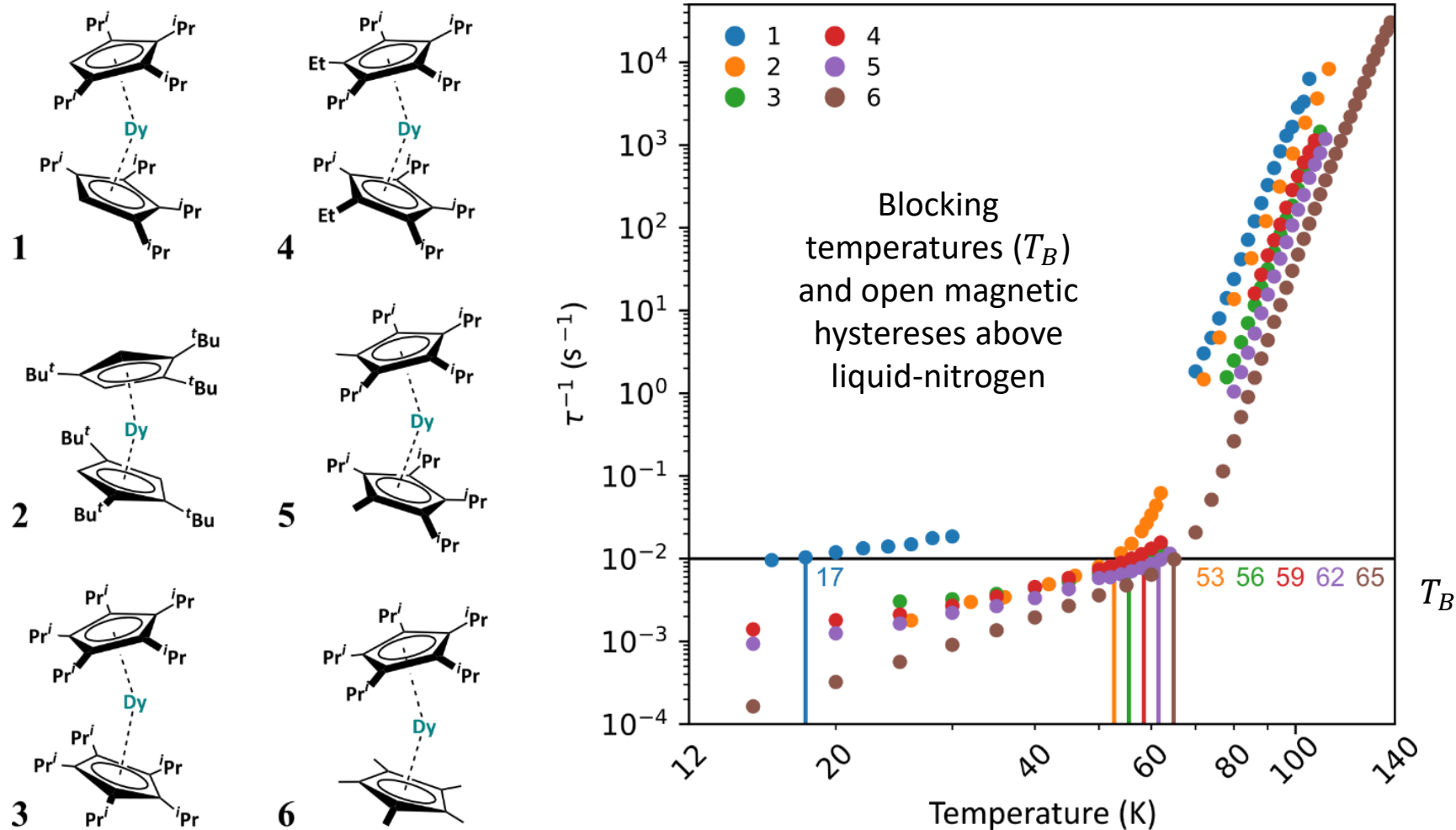
- Molecular design: from static to dynamic properties

Cyclopentadienyl-based Dy(III) are the best performing SMMs

1, 3-5: *Chem. Sci.* **2019**, *9*, 8492

2: *Nature*, **2017**, *548*, 439

6: *Science* **2018**, *362*, 6421



# Single Molecule Magnets

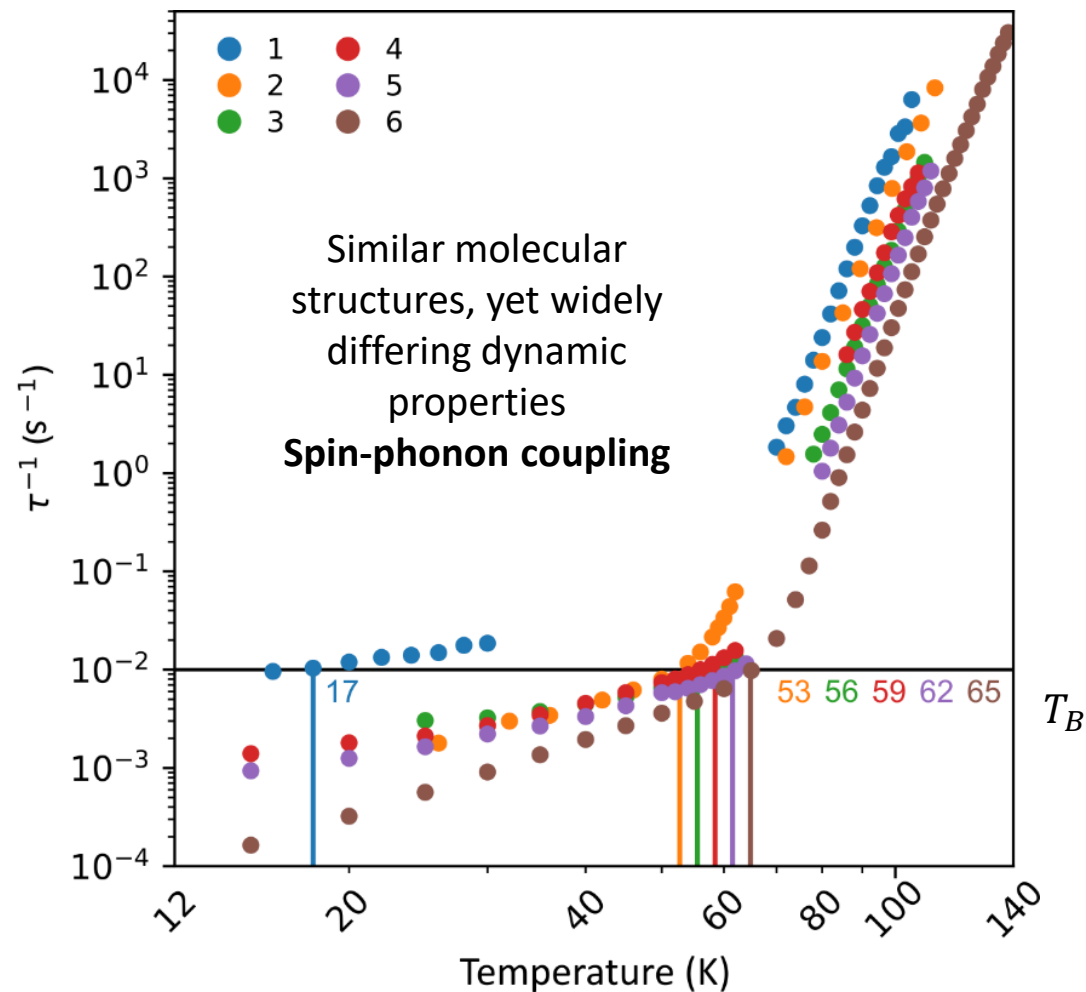
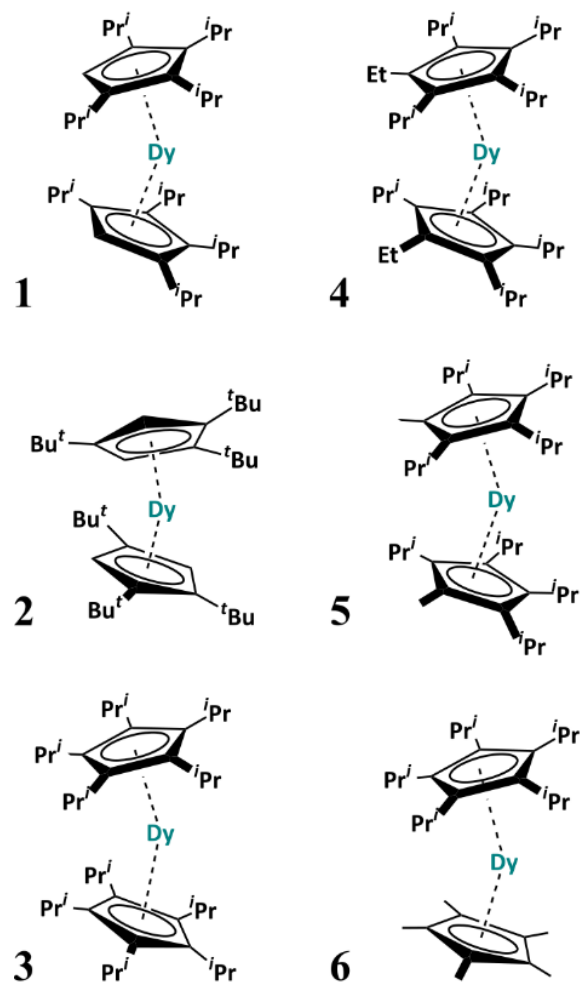
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# Single Molecule Magnets

---

- *Our approach to ab initio spin dynamics*

# Single Molecule Magnets

- *Our approach to ab initio spin dynamics*

**Geometry** optimisation & **frequency** calculation of gas-phase molecule (DFT-PBE)

**Vibronic coupling** – distort along normal modes, calculate electronic structure (CASSCF) and refer to equilibrium one (CFPs)

- Taylor series for CFPs' dependence on distortion to describe *spin-phonon coupling* Hamiltonian

$$\hat{H}_{CF} = \sum_{k=2,4,6} \sum_{q=-k}^k B_k^q \theta_k \hat{O}_k^q$$

$$B_k^q(Q_j, Q'_j, \dots) = B_k^q(Q_{eq}) + \sum_j^{3N-6} Q_j \left( \frac{\partial B_k^q}{\partial Q_j} \right)_{eq}$$

$$\hat{H}_{SP,j} = \sum_{k=2,4,6} \sum_{q=-k}^k B_k^q(Q_j) \theta_k \hat{O}_k^q = \sum_{k=2,4,6} \sum_{q=-k}^k Q_j \left( \frac{\partial B_k^q}{\partial Q_j} \right)_{eq} \theta_k \hat{O}_k^q$$

# Single Molecule Magnets

- *Our approach to ab initio spin dynamics*

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Calculate **transition rates**  $\gamma_{if}$  for the Orbach process

$$\gamma_{if} = \frac{2\pi}{\hbar} \sum_j \left| \left\langle f \left| \sum_{k=2,4,6} \sum_{q=-k}^k \left( \frac{\partial B_k^q}{\partial Q_j} \right)_{eq} \theta_k \hat{O}_k^q \right| i \right\rangle \right|^2 |\langle n_j \pm 1 | Q | n_j \rangle|^2 \rho_j(\Delta E_{fi})$$

$\langle f | \hat{H}_{SP,j}^e | i \rangle$  → Does the vibrational motion affect Dy?

# Single Molecule Magnets

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$\langle n_j \pm 1 | Q | n_j \rangle \rightarrow$  Is the vibrational mode occupied at transition energy?

# Single Molecule Magnets

- *Our approach to ab initio spin dynamics*

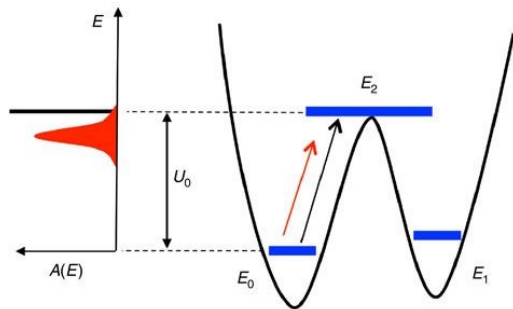
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$\rho_j(\Delta E_{fi})$  → How close in energy are the vibration to the electronic state?  
(only free parameter in our approach)



# Single Molecule Magnets

- *Our approach to ab initio spin dynamics*

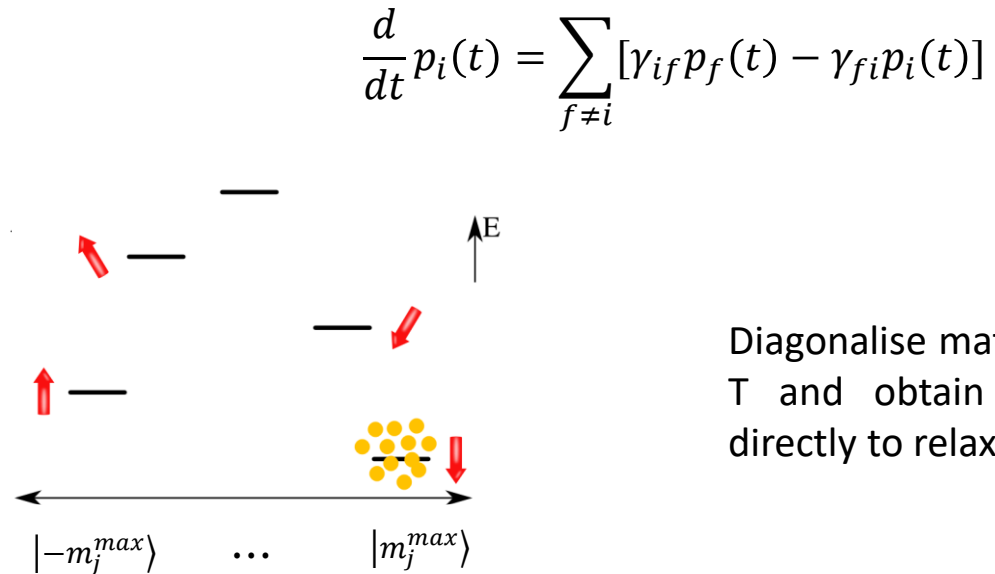
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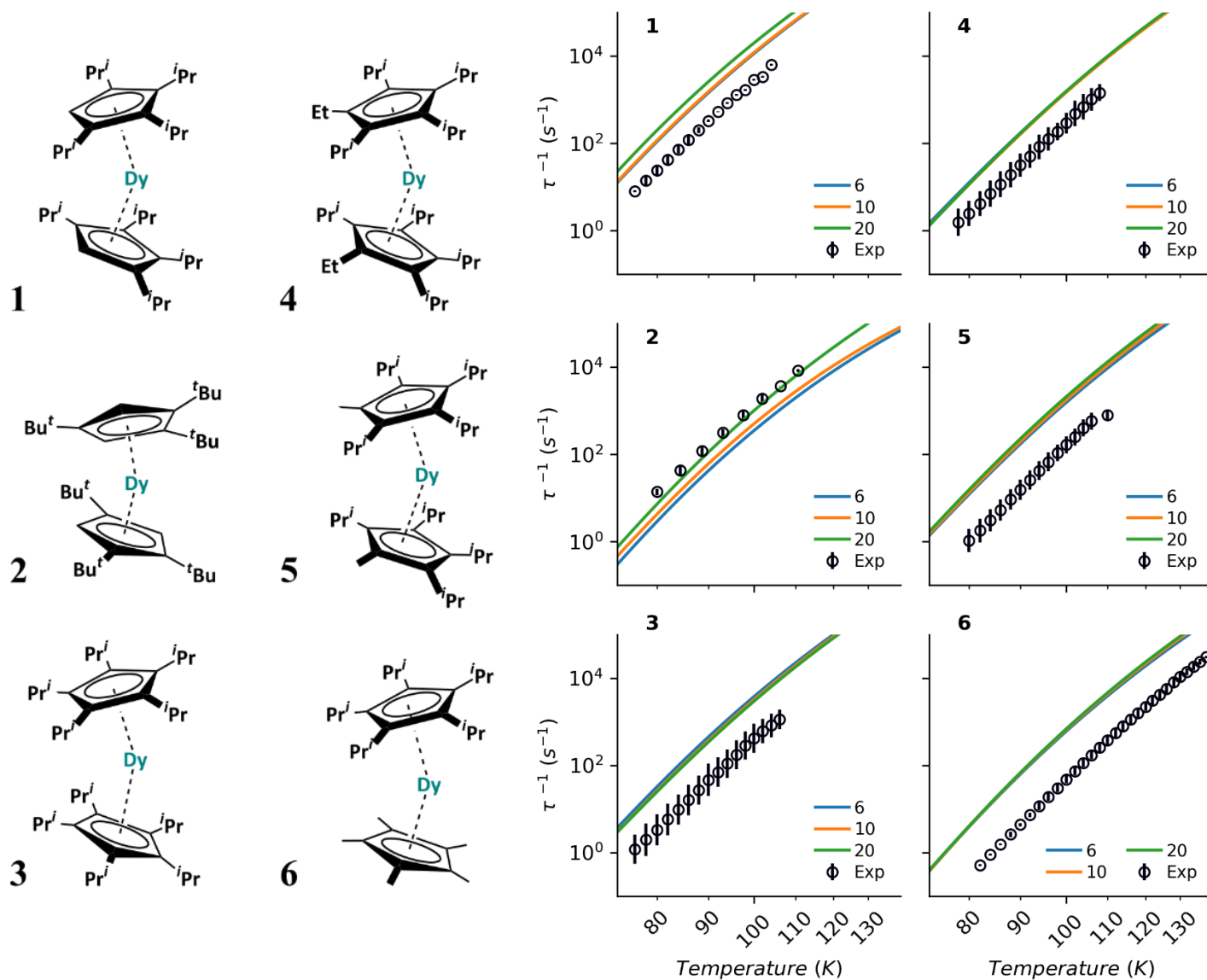
Simulate **dynamics** as a classical kinetic process master matrix



Diagonalise matrix  $\mathbf{\Gamma}$  at different  $T$  and obtain  $\tau$  to compare directly to relaxation profiles.

# Single Molecule Magnets

- Our approach to *ab initio* spin dynamics – comparison to experiment (JACS. 2021, 143, 15, 5943)

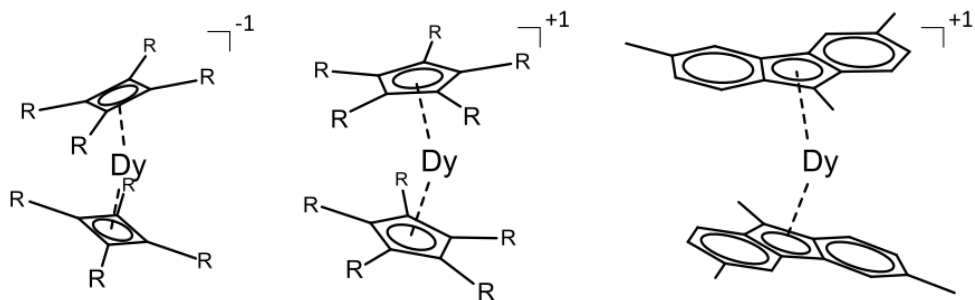


- FWHM independent.
- Consistent overestimation of  $\tau$  by a factor of *ca.* 10.
- Correct ordering of the calculated  $\tau$ .
- Our method can be confidently employed as a predictive tool.

Comparison of experimental (circles) and *ab initio* calculated (lines) relaxation rates for 1–6. Fixed fwhm line widths of 6 (blue), 10 (orange), and 20 cm<sup>-1</sup> (green) are employed.

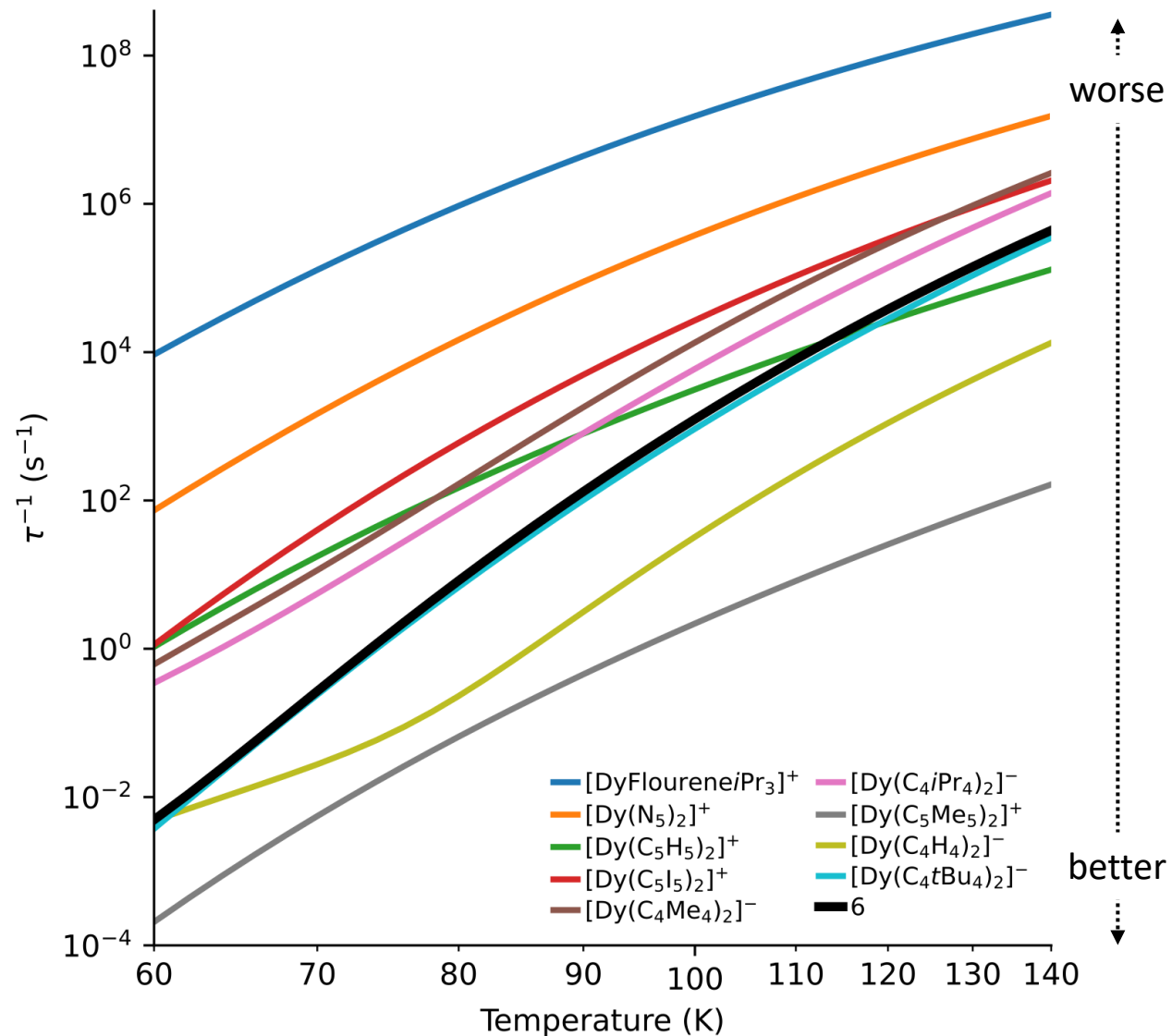
# Single Molecule Magnets

- Our approach to *ab initio* spin dynamics – can we do **better**? (*JACS.* **2021**, *143*, 15, 5943)

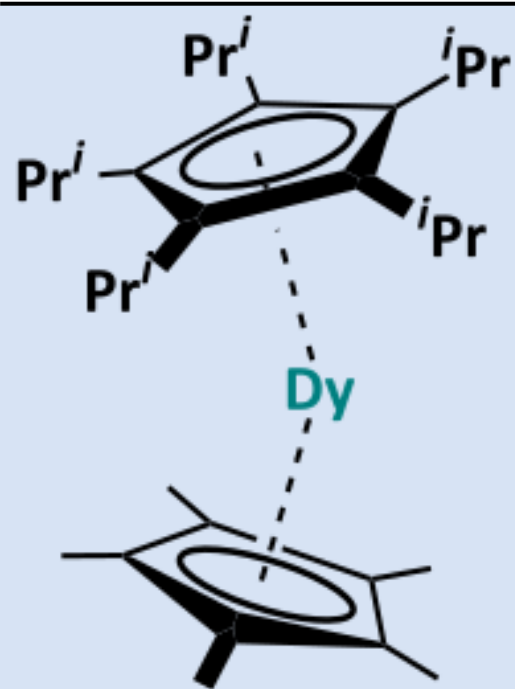


Proposed compounds aim at maximising  $U^{eff}$  by tweaking:

- Distances to ligands.
- Angle between ligands and Dy.
- Charge of ligands.



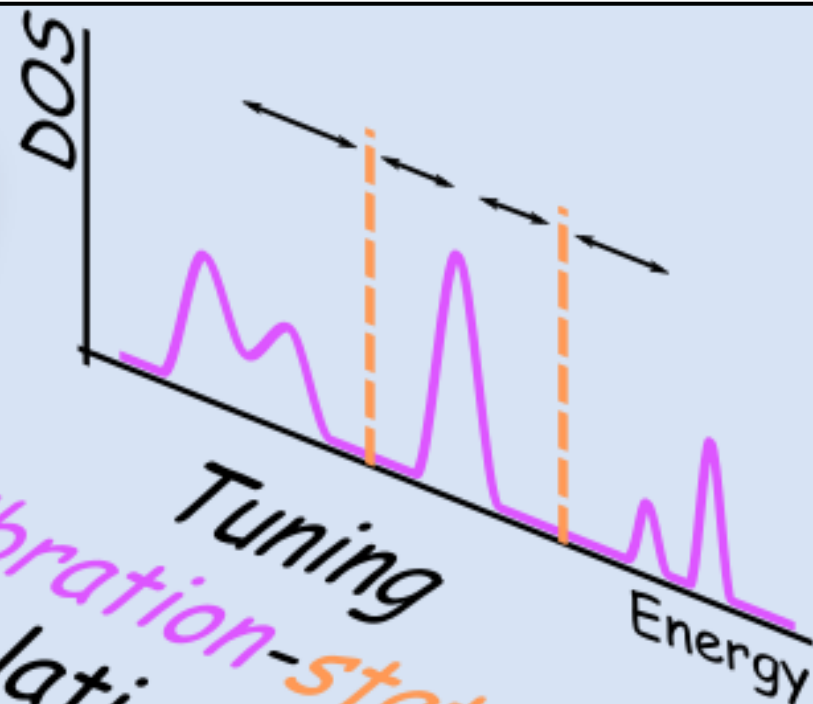
# Single Molecule Magnets



STATIC

DYNAMIC

$U_{\text{eff}}$  limit reached



SMM DESIGN STRATEGIES

# Outline

---

- Motivation
- Molecular Magnetism
  - Experimental characterisation
  - Theory and computation
- Single Molecule Magnets
  - What makes an SMMs and why do we focus on Dy(III)?
  - Molecular design: from static to dynamic properties
- Organic Radicals
  - How to trick metal-free molecules into not forming bonds: topology
  - Impact of structural flexibility
- What's next?

# Organic Radicals

---

- *How to trick metal-free molecules into not forming bonds: topological stabilisation of open-shell states*

Multiple ways to generate **monoradicals**: captodative, photolysis, redox – *not* discussed here

**Di/polyradicals** through extended  $\pi$ -systems and topological criteria: polycyclic aromatic hydrocarbons (PAHs)

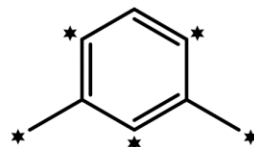
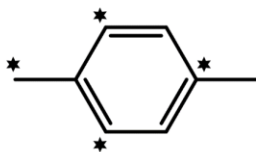
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- Lattice classification → odd alternant with non-disjoint non-bonding MOs are target



$$N = n^{starred} - n^{non-starred}$$

$$N^{even} = 0 / N^{odd} = 2$$

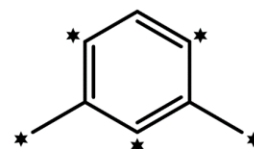
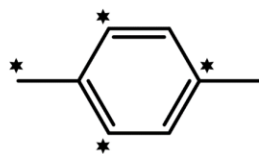
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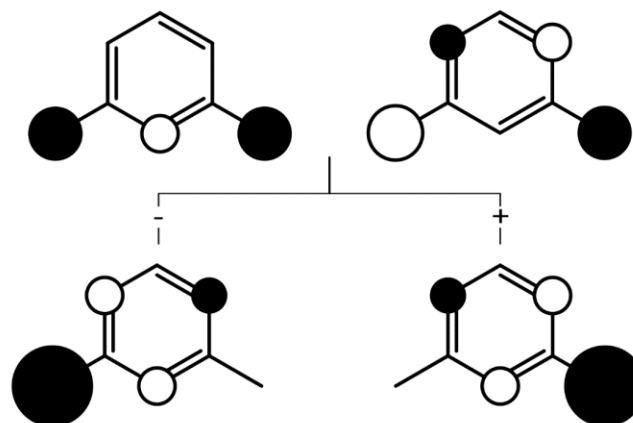
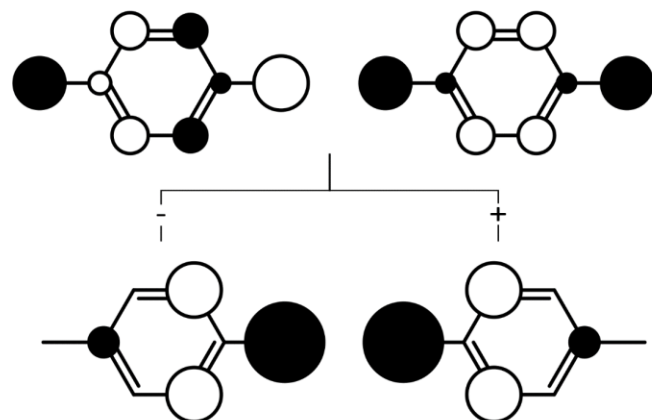
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How much do NBMOs overlap?

A little for even / A lot for odd

(because orthogonal and degenerate,  
 $S_{ab} \sim 0$  while  $K_{ab} \neq 0$  with  $\Delta E_{ST} = -2K_{ab}$ )



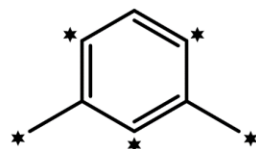
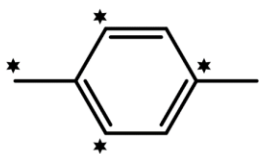
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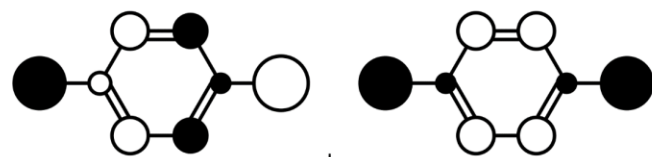
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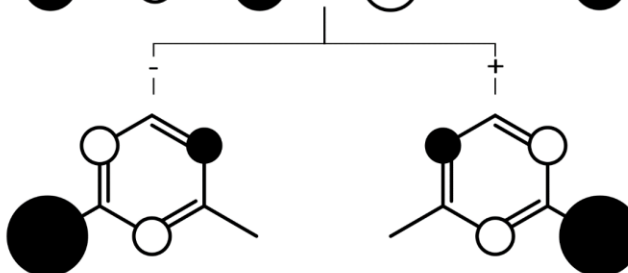
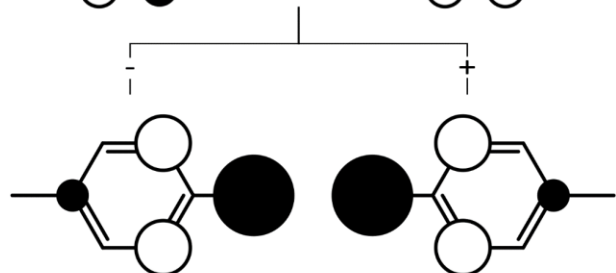
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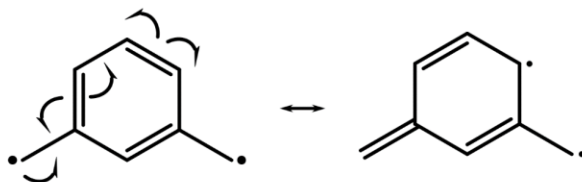
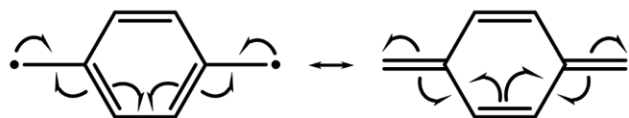
A little for even / A lot for odd



(because orthogonal and degenerate,  $S_{ab} \sim 0$  while  $K_{ab} \gg 0$  with  $\Delta E_{ST} = -2K_{ab}$ )

Broken bonds to access diamagnetic state?

0 for even / 1 for odd



# Organic Radicals

---

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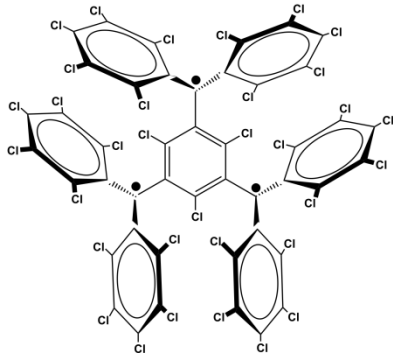
**Caution:** these are approximate guidelines derived for idealised systems. Real molecules are MUCH trickier

- Other crucial factors for radical stabilisation: kinetic (steric) and thermodynamic (delocalisation)
- Radicals are particularly sensitive to solvent effects
- Nevertheless useful

# Organic Radicals

- *Impact of structural flexibility – A molecular junction (Nano Lett. 2016, 16, 3, 2066–2071)*

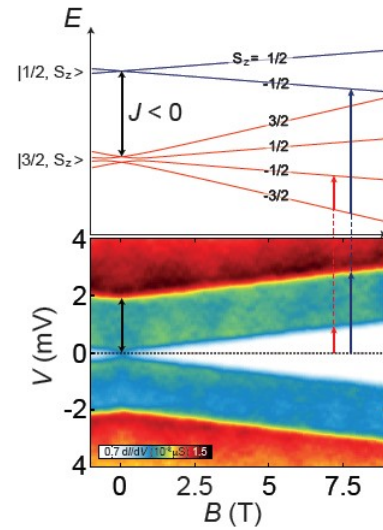
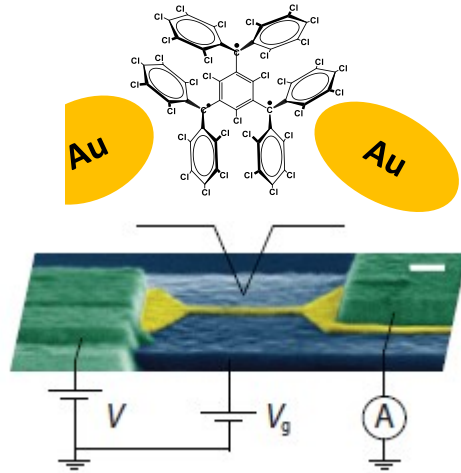
Experimental inconsistencies (IETS)



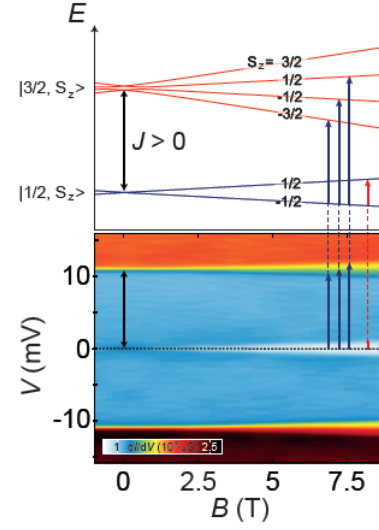
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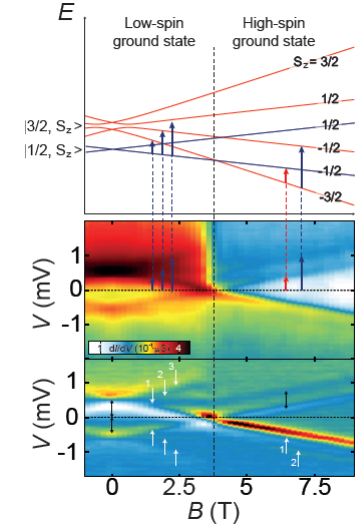
Experimental inconsistencies (IETS)



Sample #1



Sample #2



Sample #3

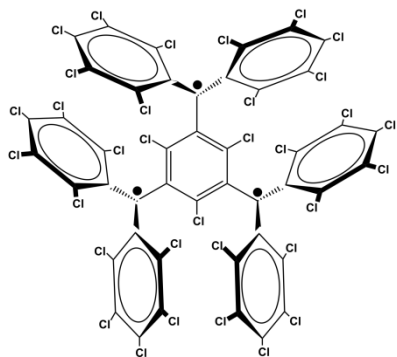
Ground state changes with sample

# Organic Radicals

- *Impact of structural flexibility – A molecular junction (Nano Lett. 2016, 16, 3, 2066–2071)*

## Computational insight

- Using Hamiltonian to calculate exchange interaction as  $\hat{H} = \sum_{\langle i,j \rangle} J_{ij} \hat{S}_i^z \hat{S}_j^z$  with  $J_{12} = J_{23} = J_{13} = J$  (isosceles)
- Broken symmetry (DFT) approach:  $HS = |\alpha\alpha\rangle$ ,  $LS = |\alpha\alpha\beta\rangle = |\alpha\beta\alpha\rangle = |\beta\alpha\alpha\rangle \rightarrow E_{|LS\rangle} - E_{|HS\rangle} = J$

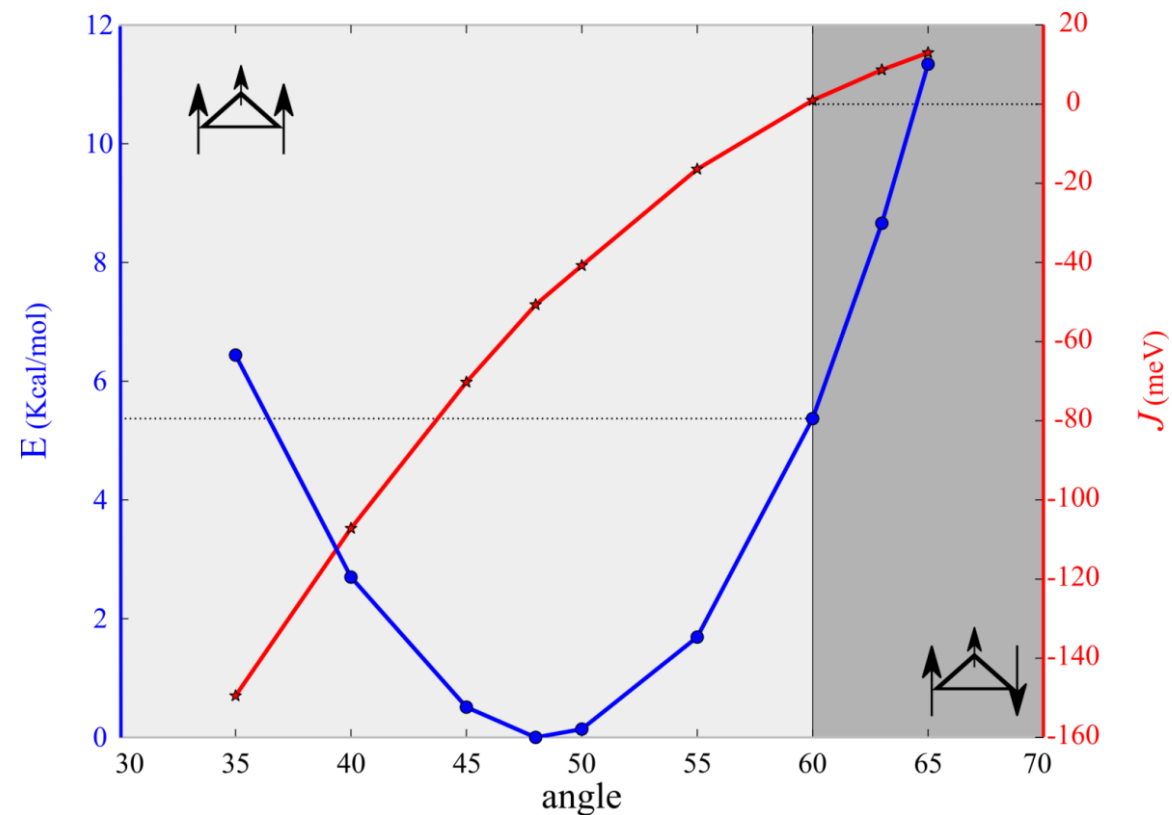
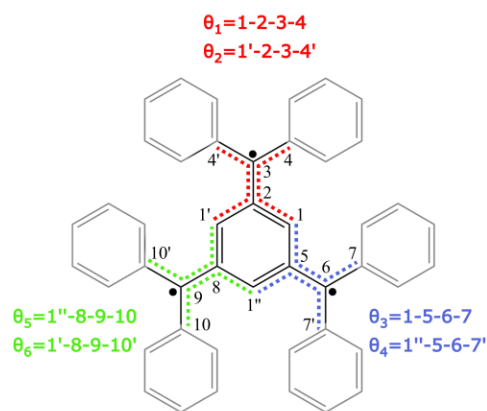


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- Explore torsion with dihedrals



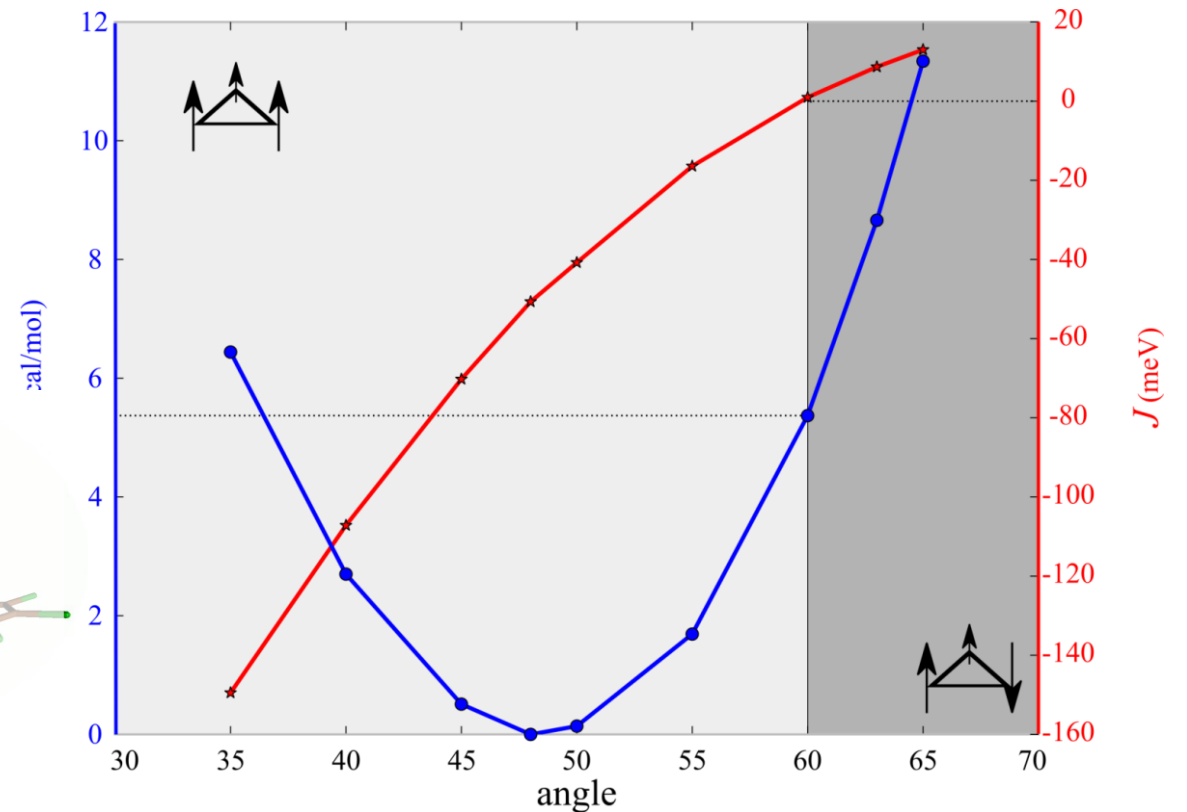
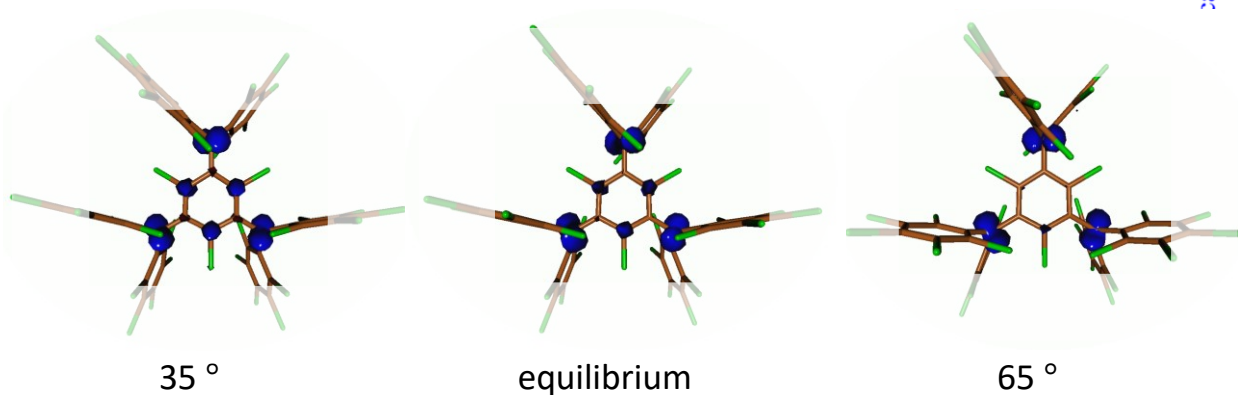
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- Explore torsion with dihedrals
- J inversion by  $\pi$ -disruption (non-disjoint  $\rightarrow$  disjoint)

Spin densities @



# Organic Radicals

---

- *Impact of structural flexibility – Chirality meets ferromagnetism (JACS, 2016, 138 (16), 5271)*

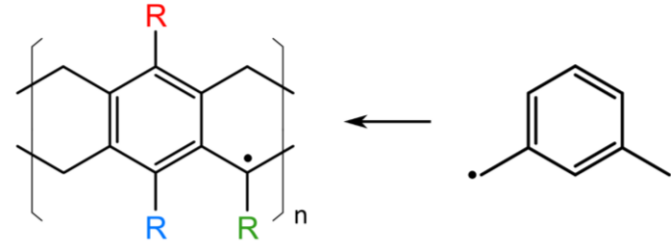


# Organic Radicals

- *Impact of structural flexibility – Chirality meets ferromagnetism (JACS, 2016, 138 (16), 5271)*

Building blocks towards low-dimensionality systems

- Previous failed synthetic efforts imposed planarity



# Organic Radicals

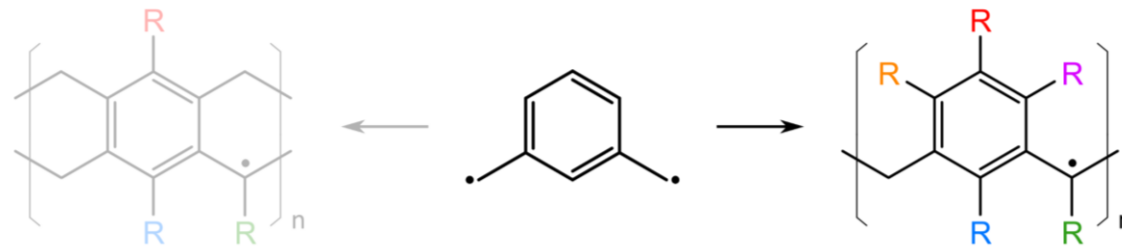
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Building blocks towards low-dimensionality systems

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Can the inherent structural flexibility be exploited?

- B3LYP calculations in periodic & gas-phase to compare structural and magnetic properties of different conformers
- Broken symmetry approach to *1<sup>st</sup> neighbours interactions*

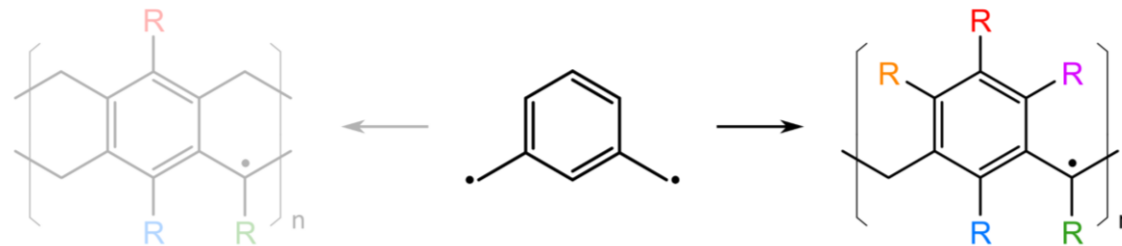


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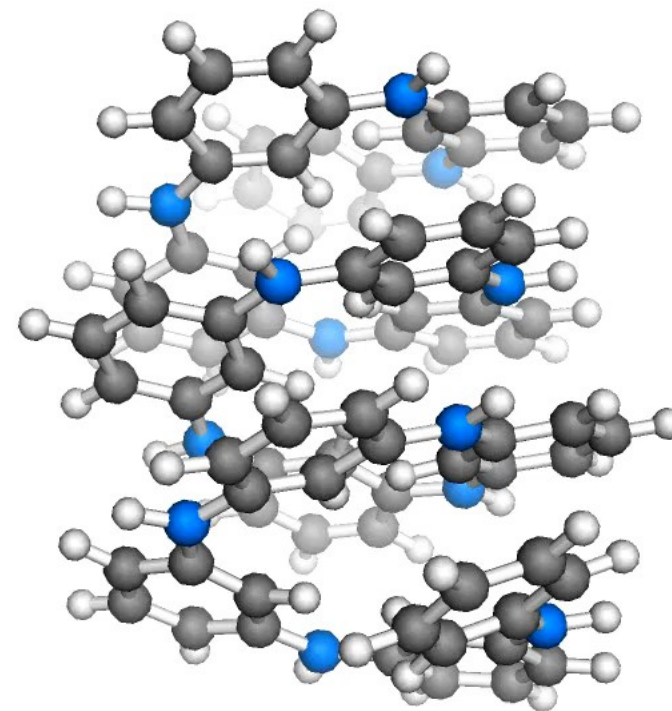
- Broken symmetry approach to *1<sup>st</sup> neighbours interactions*

- Linear vs helix:

Helix structure is **stabilised** 2.8 Kcal/mol per radical center

J is 18% **larger** in helix ( $\sim 390 \text{ cm}^{-1}$ )

- Ferromagnetism & chirality on the same molecular platform

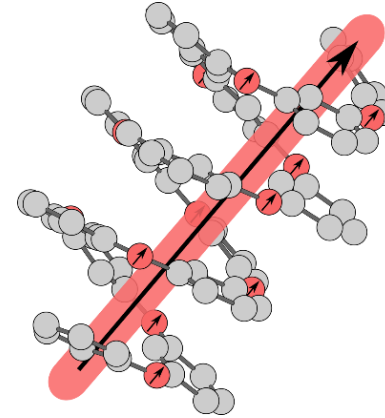


# Organic Radicals

- *Impact of structural flexibility – Chirality meets ferromagnetism*

Let's speculate

- Synthetically feasible
- Structural anisotropy induces magnetic anisotropy
- Enantiomeric separation
- Functional



# Organic Radicals

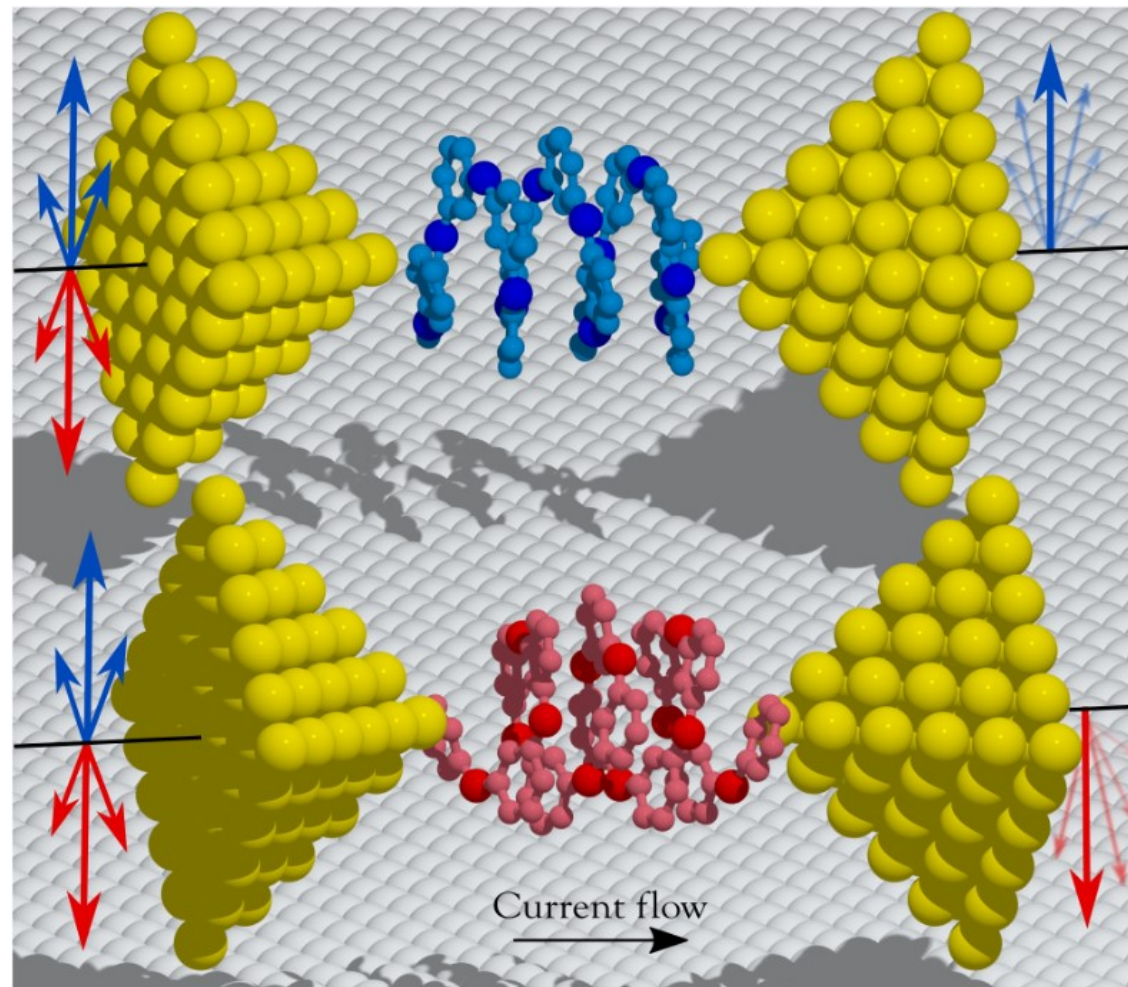
- *Impact of structural flexibility – Chirality meets ferromagnetism*

Let's speculate

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- Functional

Proposal as **spin-filter**

- Transport calculations (?)



# Acknowledgements



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